

Appendix to “On the stability of a compressible axisymmetric rotating flow
in a pipe”

By Z. Rusak & J. H. Lee

Journal of Fluid Mechanics, vol. 501 (2004), pp. 25–42

*This material has not been copy-edited or typeset by Cambridge University Press: its
format is entirely the responsibility of the author.*

Appendix A. Derivation of perturbation equations (30) and (31)

Elimination of pressure from (19) and (20) by cross differentiation in terms of x and y , respectively, followed by subtraction gives a relationship between ψ_1 , K_{1x} and ρ_1 :

$$\begin{aligned} & \left(\frac{\psi_{1xx}}{2y} + \psi_{1yy} \right)_t + w_0 \left(\frac{\psi_{1xx}}{2y} + \psi_{1yy} \right)_x - w_{0yy} \psi_{1x} + \frac{\omega K_0 \rho_0}{2y^2} K_{1x} \\ & + \gamma M_0^2 \left(\frac{\omega^2 K_0^2}{4y^2} \rho_{1x} - 2w_0 w_{0y} \rho_{1x} - w_0^2 \rho_{1xy} - \int_0^x \rho_{1ytt} dx' - 2w_{0y} \rho_{1t} - 2w_0 \rho_{1yt} \right) = 0. \end{aligned} \quad (\text{A-1})$$

Solving (A-1) for K_{1x} and substituting in linearized θ –momentum equation (21) results in:

$$\begin{aligned} K_{1t} = & \frac{\omega K_{0y}}{\rho_0} \psi_{1x} + \frac{2y^2 w_0}{\omega K_0 \rho_0} \left[\left(\frac{\psi_{1xx}}{2y} + \psi_{1yy} \right)_t + w_0 \left(\frac{\psi_{1xx}}{2y} + \psi_{1yy} \right)_x - w_{0yy} \psi_{1x} \right. \\ & \left. + \gamma M_0^2 \left(\frac{\omega^2 K_0^2}{4y^2} \rho_{1x} - 2w_0 w_{0y} \rho_{1x} - w_0^2 \rho_{1xy} - \int_0^x \rho_{1ytt} dx' - 2w_{0y} \rho_{1t} - 2w_0 \rho_{1yt} \right) \right] \end{aligned} \quad (\text{A-2})$$

Elimination of K_1 from (A-1) and (A-2) by cross differentiation in terms of t and x , respectively, followed by subtraction, and multiplying by $\omega K_0 \rho_0 / (2y^2 w_0)$, gives

$$\begin{aligned} & 2 \left(\frac{\psi_{1xx}}{2y} + \psi_{1yy} \right)_{xt} + \frac{1}{w_0} \left(\frac{\psi_{1xx}}{2y} + \psi_{1yy} \right)_{tt} + w_0 \left(\frac{\psi_{1xx}}{2y} + \psi_{1yy} \right)_{xx} \\ & + \left(\frac{\omega^2 K_0 K_{0y}}{2y^2 w_0} - w_{0yy} \right) \psi_{1xx} - \frac{w_{0yy}}{w_0} \psi_{1xt} \\ & = -\gamma M_0^2 \left[\frac{\omega^2 K_0^2}{4y^2} \left(\rho_{1xx} + \frac{\rho_{1xt}}{w_0} \right) - 4w_{0y} \rho_{1xt} - 3\rho_{1ytt} - 3w_0 \rho_{1xyt} \right. \\ & \quad \left. - \frac{1}{w_0} \int_0^x \rho_{1yttt} dx' - \frac{2w_{0y}}{w_0} \rho_{1tt} - 2w_0 w_{0y} \rho_{1xx} - w_0^2 \rho_{1xxy} \right]. \end{aligned} \quad (\text{A-3})$$

Differentiation of (A-3) with respect to x gives (30).

Differentiating (22) with respect to x gives

$$\gamma M_0^2 (\rho_0 T_{1xt} + \rho_0 w_0 T_{1xx}) = \frac{\gamma - 1}{\gamma} \left[\sqrt{2y} P_{0y} u_{1x} + \gamma M_0^2 (P_{1xt} + w_0 P_{1xx}) \right] - \rho_0 \sqrt{2y} T_{0y} u_{1x} \quad (\text{A-4})$$

From (16) we have $\rho_0 T_1 = P_1 - \rho_1 T_0$. Substituting this in (A-4) gives

$$M_0^2(P_{1xt} + w_0 P_{1xx}) - \gamma M_0^2 T_0(\rho_{1xt} + w_0 \rho_{1xx}) = \left(\frac{\gamma - 1}{\gamma} \frac{P_{0y}}{\rho_0} - T_{0y} \right) \sqrt{2y} \rho_0 u_{1x}. \quad (\text{A-5})$$

Using (20) to express P_{1xt} and P_{1xx} in (A-5) and multiplying by $-\frac{1}{w_0}$ gives

$$\begin{aligned} & \gamma M_0^2 \left[\left(\frac{T_0}{w_0} - 3M_0^2 w_0 \right) \rho_{1xt} + (T_0 - M_0^2 w_0^2) \rho_{1xx} - 3M_0^2 \rho_{1tt} - \frac{M_0^2}{w_0} \int_0^x \rho_{1ttt} dx' \right] \\ & = \left(M_0^2 w_{0y} - \frac{T_{0y}}{w_0} + \frac{\gamma - 1}{\gamma} \frac{P_{0y}}{\rho_0 w_0} \right) \psi_{1xx} - M_0^2 w_0 \psi_{1xxy} \\ & \quad - \frac{M_0^2}{w_0} \psi_{1ytt} + M_0^2 \frac{w_{0y}}{w_0} \psi_{1xt} - 2M_0^2 \psi_{1xyt}. \end{aligned} \quad (\text{A-6})$$

Differentiating (A-6) with respect to x gives (31).

Appendix B. Boundary conditions for (43) and (44)

The substitution of (42) into (32)-(41) gives boundary conditions for $\tilde{\phi}$ and $\tilde{\rho}$

$$\tilde{\phi}(x, 0) = 0, \quad \tilde{\phi}(x, 1/2) = \tilde{\phi}(0, 1/2) \quad (\text{B-1})$$

for $0 \leq x \leq x_0$ and

$$\tilde{\phi}_{xx}(0, y) = 0, \quad (\text{B-2})$$

$$\gamma M_0^2 w_0 \tilde{\rho}(0, y) = \tilde{\phi}_y(0, y), \quad (\text{B-3})$$

$$\begin{aligned} \left(\frac{T_0}{w_0} \tilde{\phi}_y(0, y) \right)_y &= \gamma M_0^2 \left(\frac{\sigma \tilde{\phi}_x(0, y)}{2y} + \omega^2 \frac{K_0^2}{4y^2 w_0} \tilde{\phi}_y(0, y) \right) \\ &\text{with } \tilde{\phi}(0, 0) = \tilde{\phi}_y(0, 0) = 0, \end{aligned} \quad (\text{B-4})$$

$$\begin{aligned} & \frac{\tilde{\phi}_{xxx}(0, y)}{2y} + \tilde{\phi}_{xyy}(0, y) + \left(\frac{\omega^2 K_0 K_{0y}}{2y^2 w_0^2} - \frac{w_{0yy}}{w_0} \right) \tilde{\phi}_x(0, y) + \sigma \frac{\tilde{\phi}_{yy}(0, y)}{w_0} \\ & = -\gamma M_0^2 \left[\left(\frac{\omega^2 K_0^2}{4y^2 w_0} - 2w_{0y} \right) \tilde{\rho}_x(0, y) - w_0 \tilde{\rho}_{xy}(0, y) - 2\sigma \frac{w_{0y}}{w_0} \tilde{\rho}(0, y) - 2\sigma \tilde{\rho}_y(0, y) \right], \end{aligned} \quad (\text{B-5})$$

$$\begin{aligned} & \sigma \frac{\tilde{\phi}_{xxx}(0, y)}{y w_0} + \sigma \frac{2\tilde{\phi}_{xyy}(0, y)}{w_0} + \frac{\sigma^2}{w_0^2} \tilde{\phi}_{yy}(0, y) + \frac{\tilde{\phi}_{xxxx}(0, y)}{2y} - \sigma \frac{w_{0yy}}{w_0^2} \tilde{\phi}_x(0, y) \\ & = -\frac{\gamma M_0^2}{w_0} \left[\frac{\omega^2 K_0^2}{4y^2} \left(\tilde{\rho}_{xx}(0, y) + \sigma \frac{\tilde{\rho}_x(0, y)}{w_0} \right) - 4\sigma w_{0y} \tilde{\rho}_x(0, y) - 3\sigma^2 \tilde{\rho}_y(0, y) \right. \\ & \quad \left. - 3\sigma w_0 \tilde{\rho}_{xy}(0, y) - \sigma^2 \frac{2w_{0y}}{w_0} \tilde{\rho}(0, y) - 2w_0 w_{0y} \tilde{\rho}_{xx}(0, y) - w_0^2 \tilde{\rho}_{xxy}(0, y) \right], \end{aligned} \quad (\text{B-6})$$

$$\begin{aligned} & \gamma \left[\sigma \left(\frac{T_0}{w_0} - 3M_0^2 w_0 \right) \tilde{\rho}_x(0, y) + (T_0 - M_0^2 w_0^2) \tilde{\rho}_{xx}(0, y) - 3\sigma^2 M_0^2 \tilde{\rho}(0, y) \right] \\ & = \sigma \frac{w_{0y}}{w_0} \tilde{\phi}_x(0, y) - \sigma^2 \frac{1}{w_0} \tilde{\phi}_y(0, y) - 2\sigma \tilde{\phi}_{xy}(0, y), \end{aligned} \quad (\text{B-7})$$

$$\tilde{\phi}_x(x_0, y) = 0, \quad \sigma \tilde{\rho}(x_0, y) + w_0 \tilde{\rho}_x(x_0, y) = 0 \quad (\text{B-8})$$

for $0 \leq y \leq 1/2$.

Appendix C. Analysis of imaginary parts of (43) and (44)

Substituting (60) into (43) and (44), collecting terms of the orders ϵ_I , σ_I , $\epsilon_I\sigma_R$, $\epsilon_I\Delta\Omega$ and neglecting terms of the orders $O(\sigma_R^2, \sigma_I^2, \sigma_R\sigma_I, \sigma_I\epsilon_R, \sigma_I\Delta\Omega)$ and higher gives

$$\begin{aligned} & \epsilon_I \left\{ \frac{\phi_{Ixx}}{2y} + \phi_{Iyy} + \left(\frac{\Omega_1 K_0 K_{0y}}{2y^2 w_0^2} - \frac{w_{0yy}}{w_0} \right) \phi_I + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0} - 2w_{0y} \right) \rho_I - w_0 \rho_{Iy} \right] \right\}_{xxx} \\ & + \sigma_I \left\{ \frac{2}{w_0} \left(\frac{\psi_{1cxx}}{2y} + \psi_{1cyy} \right) - \frac{w_{0yy}}{w_0^2} \psi_{1c} + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0^2} - 4 \frac{w_{0y}}{w_0} \right) \rho_{1c} - 3\rho_{1cy} \right] \right\}_{xx} \\ & + \epsilon_I \sigma_R \left\{ \frac{2}{w_0} \left(\frac{\phi_{Ixx}}{2y} + \phi_{Iyy} \right) - \frac{w_{0yy}}{w_0^2} \phi_I + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0^2} - 4 \frac{w_{0y}}{w_0} \right) \rho_I - 3\rho_{Iy} \right] \right\}_{xx} \\ & + \epsilon_I \Delta\Omega \left\{ \frac{K_0 K_{0y}}{2y^2 w_0^2} \phi_I + \gamma M_0^2 \frac{K_0^2}{4y^2 w_0} \rho_I \right\}_{xxx} = 0, \end{aligned} \quad (C-1)$$

$$\begin{aligned} & \epsilon_I \left\{ \rho_I - \frac{1}{\gamma M_0^2 (T_0 - M_0^2 w_0^2)} \left[\left(M_0^2 w_{0y} - \frac{T_{0y}}{w_0} + \frac{(\gamma - 1) M_0^2 \Omega_1 K_0^2}{4y^2 w_0} \right) \phi_I - M_0^2 w_0 \phi_{Iy} \right] \right\}_{xxx} \\ & + \frac{\sigma_I}{T_0 - M_0^2 w_0^2} \left\{ \frac{T_0 - 3M_0^2 w_0^2}{w_0} \rho_{1c} - \frac{w_{0y}}{\gamma w_0} \psi_{1c} + \frac{2}{\gamma} \psi_{1cy} \right\}_{xx} \\ & + \frac{\epsilon_I \sigma_R}{T_0 - M_0^2 w_0^2} \left\{ \frac{T_0 - 3M_0^2 w_0^2}{w_0} \rho_I - \frac{w_{0y}}{\gamma w_0} \phi_I + \frac{2}{\gamma} \phi_{Iy} \right\}_{xx} \\ & + \epsilon_I \Delta\Omega \left\{ \frac{\gamma - 1}{\gamma (T_0 - M_0^2 w_0^2)} \frac{K_0^2}{4y^2 w_0} \phi_I \right\}_{xxx} = 0. \end{aligned} \quad (C-2)$$

From (B-1)-(B-8), the boundary conditions for these equations are:

$$\phi_I(x, 0) = 0, \quad \phi_I(x, 1/2) = \phi_I(0, 1/2) \quad (C-3)$$

for $0 \leq x \leq x_0$ and

$$\phi_{Ixx}(0, y) = 0, \quad (C-4)$$

$$\gamma M_0^2 w_0 \rho_I(0, y) = \phi_{Iy}(0, y), \quad (C-5)$$

$$\begin{aligned} & \epsilon_I \left\{ \left(\frac{T_0}{w_0} \phi_{Iy}(0, y) \right)_y - \gamma M_0^2 \frac{\Omega_1 K_0^2}{4y^2 w_0} \phi_{Iy}(0, y) \right\} - \sigma_I \left\{ \gamma M_0^2 \frac{\psi_{1cx}(0, y)}{2y} \right\} \\ & - \epsilon_I \sigma_R \left\{ \gamma M_0^2 \frac{\phi_{Ix}(0, y)}{2y} \right\} - \epsilon_I \Delta\Omega \left\{ \gamma M_0^2 \frac{K_0^2}{4y^2 w_0} \phi_{Iy}(0, y) \right\} = 0 \\ & \text{with } \phi_I(0, 0) = \phi_{Iy}(0, 0) = 0, \end{aligned} \quad (C-6)$$

$$\begin{aligned} & \epsilon_I \left\{ \frac{\phi_{Ixxx}(0, y)}{2y} + \phi_{Ixyy}(0, y) + \left(\frac{\Omega_1 K_0 K_{0y}}{2y^2 w_0^2} - \frac{w_{0yy}}{w_0} \right) \phi_{Ix}(0, y) \right. \\ & \left. + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0} - 2w_{0y} \right) \rho_{Ix}(0, y) - w_0 \rho_{Ixy}(0, y) \right] \right\} \\ & + \epsilon_I \sigma_R \left\{ \frac{\phi_{Iyy}(0, y)}{w_0} - 2\gamma M_0^2 \left(\frac{w_{0y}}{w_0} \rho_I(0, y) + \rho_{Iy}(0, y) \right) \right\} \end{aligned}$$

$$+\epsilon_I \Delta \Omega \left\{ \frac{K_0 K_{0y}}{2y^2 w_0^2} \phi_{Ix}(0, y) + \gamma M_0^2 \frac{K_0^2}{4y^2 w_0} \rho_{Ix}(0, y) \right\} = 0, \quad (\text{C-7})$$

$$\begin{aligned} & \epsilon_I \left\{ \frac{\phi_{Ixxxx}(0, y)}{2y} + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0} - 2w_{0y} \right) \rho_{Ixx}(0, y) - w_0 \rho_{Ixy}(0, y) \right] \right\} \\ & + \sigma_I \left\{ \frac{\psi_{1cxxx}(0, y)}{y w_0} + \frac{2\psi_{1cxyy}(0, y)}{w_0} - \frac{w_{0yy}}{w_0^2} \psi_{1cx}(0, y) \right. \\ & \quad \left. + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0^2} - 4 \frac{w_{0y}}{w_0} \right) \rho_{1cx}(0, y) - 3\rho_{1cxy}(0, y) \right] \right\} \\ & + \epsilon_I \sigma_R \left\{ \frac{\phi_{Ixxx}(0, y)}{y w_0} + \frac{2\phi_{Ixyy}(0, y)}{w_0} - \frac{w_{0yy}}{w_0^2} \phi_{Ix}(0, y) \right. \\ & \quad \left. + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0^2} - 4w_{0y} \right) \rho_{Ix}(0, y) - 3\rho_{Ixy}(0, y) \right] \right\} \\ & + \epsilon_I \Delta \Omega \left\{ \gamma M_0^2 \frac{K_0^2}{4y^2 w_0} \rho_{Ixx}(0, y) \right\} = 0, \end{aligned} \quad (\text{C-8})$$

$$\begin{aligned} & \epsilon_I \rho_{Ixx}(0, y) + \frac{\sigma_I}{T_0 - M_0^2 w_0^2} \left\{ \frac{T_0 - 3M_0^2 w_0^2}{w_0} \rho_{1cx}(0, y) - \frac{w_{0y}}{\gamma w_0} \psi_{1cx}(0, y) + \frac{2}{\gamma} \psi_{1cxy}(0, y) \right\} \\ & + \frac{\epsilon_I \sigma_R}{T_0 - M_0^2 w_0^2} \left\{ \frac{T_0 - 3M_0^2 w_0^2}{w_0} \rho_{Ix}(0, y) - \frac{w_{0y}}{\gamma w_0} \phi_{Ix}(0, y) + \frac{2}{\gamma} \phi_{Ixy}(0, y) \right\} = 0, \end{aligned} \quad (\text{C-9})$$

$$\phi_{Ix}(x_0, y) = 0, \quad \sigma_I \rho_{1c}(x_0, y) + w_0 \epsilon_I \rho_{Ix}(x_0, y) + \epsilon_I \sigma_R \rho_I(x_0, y) = 0 \quad (\text{C-10})$$

for $0 \leq y \leq 1/2$.

Two integrations with respect to x of (C-1) and (C-2) and the use of boundary conditions (C-4) and (C-7)-(C-10) result in

$$\begin{aligned} & \epsilon_I \left\{ \frac{\phi_{Ixx}}{2y} + \phi_{Iyy} + \left(\frac{\Omega_1 K_0 K_{0y}}{2y^2 w_0^2} - \frac{w_{0yy}}{w_0} \right) \phi_I + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0} - 2w_{0y} \right) \rho_I - w_0 \rho_{Iy} \right] \right\}_x \\ & + \sigma_I \left\{ \frac{2}{w_0} \left(\frac{\psi_{1cxx}}{2y} + \psi_{1cyy} \right) - \frac{w_{0yy}}{w_0^2} \psi_{1c} + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0^2} - 4 \frac{w_{0y}}{w_0} \right) \rho_{1c} - 3\rho_{1cy} \right] \right\} \\ & + \epsilon_I \sigma_R \left\{ \frac{2}{w_0} \left(\frac{\phi_{Ixx}}{2y} + \phi_{Iyy} \right) - \frac{w_{0yy}}{w_0^2} \phi_I + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0^2} - 4 \frac{w_{0y}}{w_0} \right) \rho_I - 3\rho_{Iy} \right] \right\} \\ & + \epsilon_I \Delta \Omega \left\{ \frac{K_0 K_{0y}}{2y^2 w_0^2} \phi_I + \gamma M_0^2 \frac{K_0^2}{4y^2 w_0} \rho_I \right\}_x = \epsilon_I \sigma_R f_1(y), \end{aligned} \quad (\text{C-11})$$

$$\begin{aligned} & \epsilon_I \left\{ \rho_I - \frac{1}{\gamma M_0^2 (T_0 - M_0^2 w_0^2)} \left[\left(M_0^2 w_{0y} - \frac{T_{0y}}{w_0} + \frac{(\gamma - 1) M_0^2 \Omega_1 K_0^2}{4y^2 w_0} \right) \phi_I - M_0^2 w_0 \phi_{Iy} \right] \right\}_x \\ & + \frac{\sigma_I}{T_0 - M_0^2 w_0^2} \left\{ \frac{T_0 - 3M_0^2 w_0^2}{w_0} \rho_{1c} - \frac{w_{0y}}{\gamma w_0} \psi_{1c} + \frac{2}{\gamma} \psi_{1cy} \right\} \\ & + \frac{\epsilon_I \sigma_R}{T_0 - M_0^2 w_0^2} \left\{ \frac{T_0 - 3M_0^2 w_0^2}{w_0} \rho_I - \frac{w_{0y}}{\gamma w_0} \phi_I + \frac{2}{\gamma} \phi_{Iy} \right\} \\ & + \epsilon_I \Delta \Omega \left\{ \frac{\gamma - 1}{\gamma (T_0 - M_0^2 w_0^2)} \frac{K_0^2}{4y^2 w_0} \phi_I \right\}_x = \sigma_I f_2(y) + \epsilon_I \sigma_R f_3(y), \end{aligned} \quad (\text{C-12})$$

Here

$$\begin{aligned}
f_1(y) &= \frac{\phi_{Iyy}(0, y)}{w_0} - \frac{w_{0yy}}{w_0^2} \phi_I(0, y) + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0^2} - 2 \frac{w_{0y}}{w_0} \right) \rho_I(0, y) - \rho_{Iy}(0, y) \right], \\
f_2(y) &= -\frac{1}{T_0 - M_0^2 w_0^2} \left(2M_0^2 w_0 S(y) + \frac{w_{0y}}{\gamma w_0} \Phi(y) - \frac{2}{\gamma} \Phi_y(y) \right), \\
f_3(y) &= -\frac{1}{T_0 - M_0^2 w_0^2} \left(2M_0^2 w_0 \rho_I(x_0, y) + \frac{w_{0y}}{\gamma w_0} \phi_I(x_0, y) - \frac{2}{\gamma} \phi_{Iy}(x_0, y) \right).
\end{aligned} \tag{C-13}$$

Solution of (C-12) for $\epsilon_I \rho_{Ix}$, substitution of the result in (C-11), multiplication by $(T_0 - M_0^2 w_0^2)/T_0$ and an additional integration with respect to x gives (61).

Appendix D. Analysis of real parts of (43) and (44)

Substituting expansions (60) into (43) and (44) and neglecting terms of the orders $O(\sigma_R^2, \sigma_I^2, \sigma_R \epsilon_R, \sigma_I \epsilon_I, \sigma_R \Delta \Omega)$ and higher gives

$$\begin{aligned}
\epsilon_R \left\{ \frac{\phi_{Rxx}}{2y} + \phi_{Ryy} + \left(\frac{\Omega_1 K_0 K_{0y}}{2y^2 w_0^2} - \frac{w_{0yy}}{w_0} \right) \phi_R + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0} - 2w_{0y} \right) \rho_R - w_0 \rho_{Ry} \right] \right\}_{xxx} \\
+ \sigma_R \left\{ \frac{2}{w_0} \left(\frac{\psi_{1cxx}}{2y} + \psi_{1cyy} \right) - \frac{w_{0yy}}{w_0^2} \psi_{1c} + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0^2} - 4 \frac{w_{0y}}{w_0} \right) \rho_{1c} - 3\rho_{1cy} \right] \right\}_{xx} \\
+ \Delta \Omega \left\{ \frac{K_0 K_{0y}}{2y^2 w_0^2} \psi_{1c} + \gamma M_0^2 \frac{K_0^2}{4y^2 w_0} \rho_{1c} \right\}_{xxx} = 0,
\end{aligned} \tag{D-1}$$

$$\begin{aligned}
\epsilon_R \left\{ \rho_R - \frac{1}{\gamma M_0^2 (T_0 - M_0^2 w_0^2)} \left[\left(M_0^2 w_{0y} - \frac{T_{0y}}{w_0} + (\gamma - 1) M_0^2 \frac{\Omega_1 K_0^2}{4y^2 w_0} \right) \phi_R - M_0^2 w_0 \phi_{Ry} \right] \right\}_{xxx} \\
+ \frac{\sigma_R}{T_0 - M_0^2 w_0^2} \left\{ \frac{T_0 - 3M_0^2 w_0^2}{w_0} \rho_{1c} - \frac{w_{0y}}{\gamma w_0} \psi_{1c} + \frac{2}{\gamma} \psi_{1cy} \right\}_{xx} \\
- \Delta \Omega \left\{ \frac{(\gamma - 1)}{\gamma (T_0 - M_0^2 w_0^2)} \frac{K_0^2}{4y^2 w_0} \psi_{1c} \right\}_{xxx} = 0.
\end{aligned} \tag{D-2}$$

From (B-1)-(B-8), the boundary conditions for these equations are:

$$\phi_R(x, 0) = 0, \quad \phi_R(x, 1/2) = \phi_R(0, 1/2) \tag{D-3}$$

for $0 \leq x \leq x_0$ and

$$\phi_{Rxx}(0, y) = 0, \tag{D-4}$$

$$\gamma M_0^2 w_0 \rho_R(0, y) = \phi_{Ry}(0, y), \tag{D-5}$$

$$\begin{aligned}
\epsilon_R \left\{ \left(\frac{T_0}{w_0} \phi_{Ry}(0, y) \right)_y - \gamma M_0^2 \frac{\Omega_1 K_0^2}{4y^2 w_0} \phi_{Ry}(0, y) \right\} - \sigma_R \left\{ \gamma M_0^2 \frac{\psi_{1cx}(0, y)}{2y} \right\} = 0 \\
\text{with } \phi_R(0, 0) = \phi_{Ry}(0, 0) = 0,
\end{aligned} \tag{D-6}$$

$$\begin{aligned}
\epsilon_R \left\{ \frac{\phi_{Rxxx}(0, y)}{2y} + \phi_{Rxyy}(0, y) + \left(\frac{\Omega_1 K_0 K_{0y}}{2y^2 w_0^2} - \frac{w_{0yy}}{w_0} \right) \phi_{Rx}(0, y) \right. \\
\left. + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0} - 2w_{0y} \right) \rho_{Rx}(0, y) - w_0 \rho_{Rxy}(0, y) \right] \right\}
\end{aligned}$$

$$+\Delta\Omega \left\{ \frac{K_0 K_{0y}}{2y^2 w_0^2} \psi_{1cx}(0, y) + \gamma M_0^2 \frac{K_0^2}{4y^2 w_0} \rho_{1cx}(0, y) \right\} = 0, \quad (\text{D-7})$$

$$\begin{aligned} \epsilon_R \left\{ \frac{\phi_{Rxxx}(0, y)}{2y} + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0} - 2w_{0y} \right) \rho_{Rxx}(0, y) - w_0 \rho_{Rxy}(0, y) \right] \right\} \\ + \sigma_R \left\{ \frac{\psi_{1cxxx}(0, y)}{yw_0} + \frac{2\psi_{1cxyy}(0, y)}{w_0} - \frac{w_{0yy}}{w_0^2} \psi_{1cx}(0, y) \right. \\ \left. + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0^2} - 4\frac{w_{0y}}{w_0} \right) \rho_{1cx}(0, y) - 3\rho_{1cxy}(0, y) \right] \right\} = 0, \quad (\text{D-8}) \end{aligned}$$

$$\begin{aligned} \epsilon_R \rho_{Rxx}(0, y) \\ + \frac{\sigma_R}{T_0 - M_0^2 w_0^2} \left\{ \frac{T_0 - 3M_0^2 w_0^2}{w_0} \rho_{1cx}(0, y) - \frac{w_{0y}}{\gamma w_0} \psi_{1cx}(0, y) + \frac{2}{\gamma} \psi_{1cxy}(0, y) \right\} = 0, \quad (\text{D-9}) \end{aligned}$$

$$\phi_{Rx}(x_0, y) = 0, \quad \sigma_R \rho_{1c}(x_0, y) + \epsilon_R w_0 \rho_{Rx}(x_0, y) = 0 \quad (\text{D-10})$$

for $0 \leq y \leq 1/2$.

Two integrations of (D-1) and (D-2) with respect to x and the use of boundary conditions (D-4), (D-7)-(D-10) result in

$$\begin{aligned} \epsilon_R \left\{ \frac{\phi_{Rxx}}{2y} + \phi_{Ryy} + \left(\frac{\Omega_1 K_0 K_{0y}}{2y^2 w_0^2} - \frac{w_{0yy}}{w_0} \right) \phi_R + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0} - 2w_{0y} \right) \rho_R - w_0 \rho_{Ry} \right] \right\}_x \\ + \sigma_R \left\{ \frac{2}{w_0} \left(\frac{\psi_{1cxx}}{2y} + \psi_{1cyy} \right) - \frac{w_{0yy}}{w_0^2} \psi_{1c} + \gamma M_0^2 \left[\left(\frac{\Omega_1 K_0^2}{4y^2 w_0^2} - 4\frac{w_{0y}}{w_0} \right) \rho_{1c} - 3\rho_{1cy} \right] \right\} \\ + \Delta\Omega \left\{ \frac{K_0 K_{0y}}{2y^2 w_0^2} \psi_{1c} + \gamma M_0^2 \frac{K_0^2}{4y^2 w_0} \rho_{1c} \right\}_x = 0, \quad (\text{D-11}) \end{aligned}$$

$$\begin{aligned} \epsilon_R \left\{ \rho_R - \frac{1}{\gamma M_0^2 (T_0 - M_0^2 w_0^2)} \left[\left(M_0^2 w_{0y} - \frac{T_{0y}}{w_0} + (\gamma - 1) M_0^2 \frac{\Omega_1 K_0^2}{4y^2 w_0} \right) \phi_R - M_0^2 w_0 \phi_{Ry} \right] \right\}_x \\ + \frac{\sigma_R}{T_0 - M_0^2 w_0^2} \left\{ \frac{T_0 - 3M_0^2 w_0^2}{w_0} \rho_{1c} - \frac{w_{0y}}{\gamma w_0} \psi_{1c} + \frac{2}{\gamma} \psi_{1cy} \right\} \\ - \Delta\Omega \left\{ \frac{(\gamma - 1)}{\gamma (T_0 - M_0^2 w_0^2)} \frac{K_0^2}{4y^2 w_0} \psi_{1c} \right\}_x = \sigma_R f_2(y). \quad (\text{D-12}) \end{aligned}$$

Here $f_2(y)$ is defined in (C-13). Also, the conditions in (D-5) and (D-6) can be solved and show that

$$\begin{aligned} \epsilon_R \phi_R(0, y) &= \sigma_R \gamma M_0^2 \frac{\pi}{4x_0} \int_0^y \exp(\alpha(y')) \left[\int_0^{y'} g(y'') \exp(-\alpha(y'')) dy'' \right] dy', \\ \gamma M_0^2 w_0 \rho_R(0, y) &= \sigma_R \gamma M_0^2 \frac{\pi}{4x_0} \exp(\alpha(y)) \left[\int_0^{y'} g(y'') \exp(-\alpha(y'')) dy'' \right] \quad (\text{D-13}) \end{aligned}$$

where

$$\alpha(y) = - \int_0^y p(y') dy', \quad p(y) = \frac{w_0}{T_0} \left(\frac{T_0}{w_0} \right)_y - \gamma M_0^2 \Omega_1 \frac{K_0^2}{4y^2 T_0}, \quad g(y) = \frac{\Phi(y) w_0(y)}{y T_0(y)}.$$

Note that $\phi_R(0, y) = 0$ when $M_0 = 0$. Also, note that in the general case $\phi_R(0, 1/2)$ is now determined and may not be zero. For example, in the case of a solid-body rotation

profile where $K_0 = 2y$ and $w_0 = T_0 = 1$ we find $p(y) = -\gamma M_0^2 \Omega_1$, $\alpha(y) = \gamma M_0^2 \Omega_1 y$, and then

$$\phi_R(0, y) = \gamma M_0^2 \frac{\pi}{4x_0} \int_0^y \exp(\gamma M_0^2 \Omega_1 y) \left[\int_0^{y'} \frac{\Phi(y)}{y} \exp(-\gamma M_0^2 \Omega_1 y'') dy'' \right] dy'. \quad (\text{D-14})$$

Examples of calculating $\phi_R(0, y)$ according to (D-14) are shown in Fig. D-1 for various Mach numbers. It is clear that $\phi_R(0, 1/2)$ is not zero.

Solving (D-12) for $\epsilon_R \rho_{Rx}$ and substituting in (D-11), multiplying by $(T_0 - M_0^2 w_0^2)/T_0$, and integrating again with respect to x gives (63).

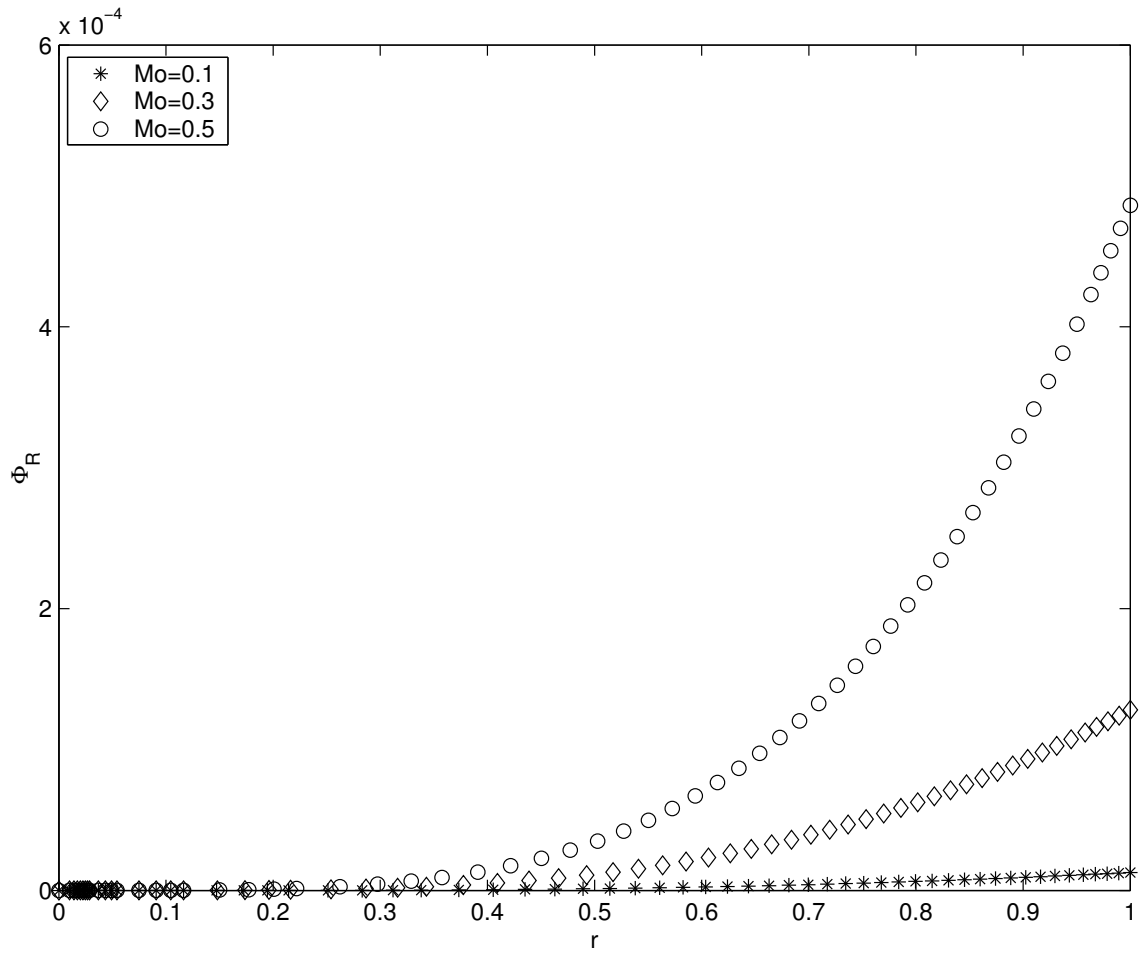


Figure 1: D-1: Solutions of $\phi_R(0, y)$ for a solid-body rotation flow at various Mach numbers.