## MARKOV CHAIN APPROXIMATION OF

ONE-DIMENSIONAL STICKY DIFFUSIONS:

## SUPPLEMENTARY MATERIAL

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## B. Additional proofs

Proof of Theorem 1: Suppose that the unique weak solution to (2.1) and (2.2) for $\rho=\infty$ is given by $\left(X^{1}, B^{1}\right)$. Let

$$
\begin{aligned}
\phi_{t} & =t+\frac{1}{2 \rho} L_{t}^{l}\left(X^{1}\right), \\
T_{t} & =\phi_{t}^{-1}, \\
X_{t} & =X_{T_{t}}^{1}, \\
B_{t} & =B_{T_{t}}^{1}+\int_{0}^{t} I\left(X_{s}=l\right) d B_{s}^{0},
\end{aligned}
$$

where $B^{0}$ is a Brownian motion, defined on an extended probability space if needed, that is independent of $B^{1}$. The local time process is continuous and non-decreasing; hence $\phi_{t}$ is strictly increasing and continuous (see [1, Chapter II.13]). This implies that $T_{t}$ is also strictly increasing and continuous. Then $B_{t}$

[^0]is a continuous local martingale and
$$
\langle B\rangle_{t}=\left\langle B^{1}\right\rangle_{T_{t}}+\int_{0}^{t} I\left(X_{s}=l\right) d s=T_{t}+\int_{0}^{t} I\left(X_{s}=l\right) d s
$$

It also holds that

$$
\begin{aligned}
T_{t} & =\int_{0}^{T_{t}} d s-\int_{0}^{T_{t}} I\left(X_{s}^{1}=0\right) d s=\int_{0}^{T_{t}} I\left(X_{s}^{1}>l\right) d s \\
& =\int_{0}^{T_{t}} I\left(X_{s}^{1}>l\right)\left(d s+\frac{1}{2 \rho} d L_{s}^{l}\left(X^{1}\right)\right) \\
& =\int_{0}^{t} I\left(X_{T_{s}}^{1}>l\right) d \phi_{T_{s}}=\int_{0}^{t} I\left(X_{s}>l\right) d s
\end{aligned}
$$

where we use that $d L_{s}^{l}(X)$ only increases for $X_{s}=l$, and we apply the change-of-variable formula. Therefore, $\langle B\rangle_{t}=t$, and by Lévy's characterization, $B$ is a standard Brownian motion. Moreover,

$$
\begin{aligned}
X_{t}= & X_{T_{t}}^{1} \\
= & \int_{0}^{T_{t}} \mu\left(X_{s}^{1}\right) I\left(X_{s}^{1}>l\right) d s+\int_{0}^{T_{t}} \sigma\left(X_{s}^{1}\right) I\left(X_{s}^{1}>l\right) d B_{s}^{1}+\frac{1}{2} \int_{0}^{T_{t}} d L_{s}^{l}\left(X^{1}\right) \\
= & \int_{0}^{T_{t}} \mu\left(X_{s}^{1}\right) I\left(X_{s}^{1}>l\right)\left(d s+\frac{1}{2 \rho} d L_{s}^{l}\left(X^{1}\right)\right) \\
& \quad+\int_{0}^{t} \sigma\left(X_{T_{s}}^{1}\right) I\left(X_{T_{s}}^{1}>l\right) d B_{T_{s}}^{1}+\frac{1}{2} L_{T_{t}}^{l}\left(X^{1}\right) \\
& =\int_{0}^{t} \mu\left(X_{T_{s}}^{1}\right) I\left(X_{T_{s}}^{1}>l\right) d \phi_{T_{s}}+\int_{0}^{t} \sigma\left(X_{s}\right) I\left(X_{s}>l\right) d B_{s}+\frac{1}{2} L_{T_{t}}^{l}\left(X^{1}\right) \\
= & \int_{0}^{t} \mu\left(X_{s}\right) I\left(X_{s}>l\right) d s+\int_{0}^{t} \sigma\left(X_{s}\right) I\left(X_{s}>l\right) d B_{s}+\frac{1}{2} L_{t}^{l}(X)
\end{aligned}
$$

because $X_{s}=X_{T_{s}}^{1}$,

$$
d \phi_{T_{s}}=d s+\frac{1}{2 \rho} d L_{s}^{l}\left(X^{1}\right)=d s
$$

as $X^{1}$ is the unique weak solution of the reflecting case, and $L_{T_{t}}^{l}\left(X^{1}\right)=L_{t}^{l}(X)$.
This shows that $(X, B)$ solves (2.1). Furthermore,

$$
\begin{aligned}
\int_{0}^{t} I\left(X_{s}=l\right) d s & =\int_{0}^{t} I\left(X_{T_{s}}^{1}=l\right) d \phi_{T_{s}}=\int_{0}^{T_{t}} I\left(X_{s}^{1}=l\right) d \phi_{s} \\
& =\int_{0}^{T_{t}} I\left(X_{s}^{1}=l\right)\left(d s+\frac{1}{2 \rho} d L_{s}^{l}\left(X^{1}\right)\right)=\frac{1}{2 \rho} L_{t}^{l}(X)
\end{aligned}
$$

where the first term vanishes because for $X^{1}$ it holds that $I\left(X_{s}^{1}=l\right) d s=0$. The continuity of $X$ follows from the continuity of $X^{1}$ and $T$. Hence, $(X, B)$ also solves (2.2).

The next step is to show the uniqueness in law of the solution $X$. We reset the notation and suppose that $(X, B)$ solves (2.1) and (2.2). Define

$$
T_{t}=\int_{0}^{t} I\left(X_{s}>l\right) d s
$$

for $t \geq 0$. Then $T_{t}$ is continuous and strictly increasing almost surely. This can be shown by contradiction. Assume $T_{t}$ is not strictly increasing; then there exists a set

$$
\Gamma=\left\{\omega \in \Omega: T_{t_{1}}=T_{t_{2}} \text { for some } 0<t_{1}<t_{2}\right\},
$$

with $\mathbb{P}(\Gamma)>0$ and $t_{1}, t_{2}$ depending on $\omega$. Now $T_{t_{1}}=T_{t_{2}}$ implies that the process stays at the boundary for all $s \in\left[t_{1}, t_{2}\right]$, and so

$$
\Gamma \subset\left\{\omega: \int_{t_{1}}^{t_{2}} d L_{s}^{l}(X)=L_{t_{2}}^{l}(X)-L_{t_{1}}^{l}(X)>0 \text { for some } 0<t_{1}<t_{2}\right\} ;
$$

i.e., the local time increases between $t_{1}$ and $t_{2}$. On this set, $I\left(X_{s}>l\right)=0$ for all $s \in\left[t_{1}, t_{2}\right]$, and hence

$$
\Gamma \subset\left\{\omega: X_{t_{2}}=X_{t_{1}}+L_{t_{2}}^{l}(X)-L_{t_{1}}^{l}(X)>X_{t_{1}} \text { for some } 0<t_{1}<t_{2}\right\}
$$

as the drift and volatility vanish. This is a contradiction to $I\left(X_{s}>l\right)=0$, and so, in summary, $T_{t}$ is strictly increasing almost surely. The inverse of $T_{t}$, given by

$$
\phi_{t}=\inf \left\{s \geq 0: T_{s}>t\right\}
$$

is therefore also continuous and almost surely finite. As $X$ and $\phi$ are continuous, it follows that $X$ is constant on every interval $\left[\phi_{t-}, \phi_{t}\right.$ ], and so $\phi$ is in synchronization with $X$ (see [4, Definition 10.13], which refers to this as adaptedness of $X$ to the time change $\phi$ ).

Now set $X_{t}^{1}=X_{\phi_{t}}$. Then $X^{1}$ is a continuous semimartingale (see Corollary
10.12 and Lemma 10.15 in [4]), and by (2.2),

$$
\begin{aligned}
t & =T_{\phi_{t}}=\int_{0}^{\phi_{t}} I\left(X_{s}>l\right) d s=\phi_{t}-\int_{0}^{\phi_{t}} I\left(X_{s}=l\right) d s \\
& =\phi_{t}-\frac{1}{2 \rho} L_{\phi_{t}}^{l}(X)=\phi_{t}-\frac{1}{2 \rho} L_{t}^{l}\left(X^{1}\right),
\end{aligned}
$$

so

$$
\phi_{t}=t+\frac{1}{2 \rho} L_{t}^{l}\left(X^{1}\right),
$$

which shows that $\phi$ is also strictly increasing. Let $B_{t}^{1}=\int_{0}^{\phi_{t}} I\left(X_{s}>l\right) d B_{s}$. Then $B_{t}^{1}$ is a continuous local martingale with

$$
\left\langle B^{1}\right\rangle_{t}=\int_{0}^{\phi_{t}} I\left(X_{s}>l\right) d s=T_{\phi_{t}}=t
$$

and hence $B^{1}$ is a Brownian motion by Lévy's criterion. Furthermore, by (2.1) it follows that

$$
\begin{aligned}
X_{t}^{1} & =X_{0}+\int_{0}^{\phi_{t}} \mu\left(X_{s}\right) I\left(X_{s}>l\right) d s+\int_{0}^{\phi_{t}} \sigma\left(X_{s}\right) I\left(X_{s}>l\right) d B_{s}+\frac{1}{2} L_{\phi_{t}}^{l}(X) \\
& =X_{0}+\int_{0}^{t} \mu\left(X_{s}^{1}\right) I\left(X_{s}^{1}>l\right) d \phi_{s}+\int_{0}^{t} \sigma\left(X_{s}^{1}\right) I\left(X_{s}^{1}>l\right) d B_{s}^{1}+\frac{1}{2} L_{t}^{l}\left(X^{1}\right),
\end{aligned}
$$

by the change of variables formula, and

$$
\begin{aligned}
d X_{t}^{1} & =\mu\left(X_{t}^{1}\right) I\left(X_{t}^{1}>l\right) d \phi_{t}+\sigma\left(X_{t}^{1}\right) I\left(X_{t}^{1}>l\right) d B_{t}^{1}+\frac{1}{2} d L_{t}^{l}\left(X^{1}\right) \\
& =\mu\left(X_{t}^{1}\right) I\left(X_{t}^{1}>l\right)\left(d t+\frac{1}{2 \rho} d L_{t}^{l}\left(X^{1}\right)\right)+\sigma\left(X_{t}^{1}\right) I\left(X_{t}^{1}>l\right) d B_{t}^{1}+\frac{1}{2} d L_{t}^{l}\left(X^{1}\right) \\
& =\mu\left(X_{t}^{1}\right) I\left(X_{t}^{1}>l\right) d t+\sigma\left(X_{t}^{1}\right) I\left(X_{t}^{1}>l\right) d B_{t}^{1}+\frac{1}{2} d L_{t}^{l}\left(X^{1}\right) .
\end{aligned}
$$

Moreover, we have

$$
\int_{0}^{t} I\left(X_{s}^{1}=l\right) d s=\int_{0}^{t} I\left(X_{\phi_{s}}=l\right) d T_{\phi_{s}}=\int_{0}^{\phi_{t}} I\left(X_{s}=l\right) d T_{s}=0 .
$$

The last two equations show that $\left(X^{1}, B^{1}\right)$ is a unique weak solution to the system of SDEs (2.1) and (2.2) for $\rho=\infty$. Since $X^{1}$ is the unique solution to the reflecting SDE and $X_{t}=X_{\phi_{T_{t}}}=X_{T_{t}}^{1}$, the law of $X$ is also unique. Theorem 3.1 in [2] states that uniqueness in law for $X$ implies joint uniqueness in law for $(X, B)$.

We will restate some results from [8] and [9] without proof but with adjustments to incorporate the sticky boundary behavior at the left boundary.

Lemma 1. For any $f, g: \mathbb{S}_{n} \rightarrow \mathbb{R}$ with $g\left(x_{n+1}\right)=0$, we have

$$
\begin{equation*}
\sum_{x \in \mathbb{S}_{n}^{\circ}} g(x) \delta^{-} x \nabla^{-} f(x)=-\sum_{x \in \mathbb{S}_{n}^{-}} f(x) \delta^{+} x \nabla^{+} g(x)-g\left(x_{0}\right) f\left(x_{0}\right) . \tag{B.1}
\end{equation*}
$$

Under Assumption 2, it can be seen that

$$
\begin{equation*}
\sup _{x, y \in(l, r)}\left|\frac{\partial^{i}}{\partial x^{i}} \frac{\partial^{j}}{\partial y^{j}} p(t, x, y)\right|<\infty \tag{B.2}
\end{equation*}
$$

for $i, j=0,1,2$ still holds, because of results from Sturm-Liouville theory and the proof of [8, Lemma 2]. The use of (B.2) is to prove claims of the form $|g(x)| \leq C h_{n}^{\beta}$ for $\beta=0,1,2$ such that the constant $C>0$ is independent of $x$ and $n$. The application of this result will not be mentioned explicitly below.

Corollary 1. Under Assumption 2, for $n$ sufficiently large, there exist constants $C_{1}, C_{2}>0$, independent of $n$ and $x \in \mathbb{S}_{n}, y \in \mathbb{S}_{n}^{\circ}$, such that

$$
C_{1} \leq s_{n}(x) \leq C_{2}, \quad C_{1} \leq m_{n}(y) \leq C_{2}, \quad C_{1} \leq M_{n}\left(x_{0}\right) \leq C_{2} .
$$

Lemma 2. Under Assumptions 2 and 3, there exists a constant $C>0$, independent of $k$ and $n$, such that for $h_{n} \in(0, \delta)$, where $\delta$ is small enough, the following holds:

$$
\begin{equation*}
\lambda_{k}^{n} \leq C k^{2} . \tag{B.3}
\end{equation*}
$$

Proof. Let the matrix $\mathbb{M}_{n}$ again be a diagonal matrix with entries $\mathbb{M}_{n, i, i}=$ $M_{n}\left(x_{i}\right)$ for $i=0, \ldots, n$. Calculation of $\mathbb{M}_{n} \mathbb{G}_{n}$ and the choice of $M_{n}(x)$ as stated in Section 3.3 implies that

$$
\begin{aligned}
\frac{\rho}{\delta^{+} x_{0}} \beta M_{n}\left(x_{0}\right) & =\frac{-\mu\left(x_{1}\right) \delta^{+} x_{1}+\sigma^{2}\left(x_{1}\right)}{2 \delta^{-} x_{1} \delta x_{1}} M_{n}\left(x_{1}\right), \\
\frac{\mu(x) \delta^{-} x+\sigma^{2}(x)}{2 \delta^{+} x \delta x} M_{n}(x) & =\frac{-\mu\left(x^{+}\right) \delta^{+} x^{+}+\sigma^{2}\left(x^{+}\right)}{2 \delta^{-} x^{+} \delta x^{+}} M_{n}\left(x^{+}\right)
\end{aligned}
$$

for $x=x_{1}, \ldots, x_{n-1}$. Hence $\mathbb{M}_{n} \mathbb{G}_{n}$ is symmetric, and therefore $\mathbb{M}_{n}^{1 / 2} \mathbb{G}_{n} \mathbb{M}_{n}^{-1 / 2}$ is also symmetric. Furthermore, it can be seen that $\mathbb{M}_{n}^{1 / 2} \mathbb{G}_{n} \mathbb{M}_{n}^{-1 / 2}$ is similar to $\mathbb{G}_{n}$, and so both matrices have the same eigenvalues. The min-max principle derived in [9, Section 3] shows that

$$
\begin{equation*}
\lambda_{k}^{n}=\min _{U_{k}} \max _{f \in U_{k}} \frac{-f^{T} \mathbb{M}_{n}^{1 / 2} \mathbb{G}_{n} \mathbb{M}_{n}^{-1 / 2} f}{f^{T} f}=\min _{U_{k}} \max _{f \in U_{k}} \frac{\left(f,-G_{n} f\right)_{n}}{(f, f)_{n}}, \tag{B.4}
\end{equation*}
$$

where $U_{k}$ denotes a $k$-dimensional subspace of functions defined on $\mathbb{S}_{n}$ with boundary condition $f\left(x_{n+1}\right)=0$.

The upper boundary for $\lambda_{k}^{n}$ can be derived in the following way. First,

$$
\begin{aligned}
\left(f,-G_{n} f\right)_{n}= & -\sum_{x \in \mathbb{S}_{n}^{-}} f(x)\left(G_{n} f(x)\right) M_{n}(x) \\
= & -\sum_{x \in \mathbb{S}_{n}^{\circ}} f(x) \frac{1}{m_{n}(x)} \frac{\delta^{-} x}{\delta x} \nabla^{-}\left(\frac{1}{s_{n}(x)} \nabla^{+} f(x)\right) m_{n}(x) \delta x \\
& +\sum_{x \in \mathbb{S}_{n}^{-}} k(x) f(x)^{2} M_{n}(x)-\rho \beta f\left(x_{0}\right) \nabla^{+} f\left(x_{0}\right) M_{n}\left(x_{0}\right) \\
= & -\sum_{x \in \mathbb{S}_{n}^{\circ}} f(x) \delta^{-} x \nabla^{-}\left(\frac{1}{s_{n}(x)} \nabla^{+} f(x)\right)+\sum_{x \in \mathbb{S}_{n}^{-}} k(x) f(x)^{2} M_{n}(x) \\
& -\rho \beta f\left(x_{0}\right) \nabla^{+} f\left(x_{0}\right) M_{n}\left(x_{0}\right) \\
= & \sum_{x \in \mathbb{S}_{n}^{s}} \frac{\delta^{+} x}{s_{n}(x)}\left(\nabla^{+} f(x)\right)^{2}+\sum_{x \in \mathbb{S}_{n}^{-}} k(x) f(x)^{2} M_{n}(x)+\frac{1-\beta}{s_{n}\left(x_{0}\right)} f\left(x_{0}\right) \nabla^{+} f\left(x_{0}\right),
\end{aligned}
$$

using Lemma 1 and $1 / s_{n}\left(x_{0}\right)=\rho M_{n}\left(x_{0}\right)$. Furthermore, by noting that for

Scheme 1, $1-\beta=0$, and for Scheme 2, $1-\beta=O\left(h_{n}\right)$, we derive

$$
\begin{aligned}
& \sum_{x \in \mathbb{S}_{n}^{-}} \quad \frac{\delta^{+} x}{s_{n}(x)}\left(\nabla^{+} f(x)\right)^{2}+\frac{1-\beta}{s_{n}\left(x_{0}\right)} f\left(x_{0}\right) \nabla^{+} f\left(x_{0}\right) \\
& \quad \leq \frac{C_{1}}{h_{n}} \sum_{x \in \mathbb{S}_{n}^{-}}\left(f\left(x^{+}\right)-f(x)\right)^{2}+\frac{C_{2}}{h_{n}}|1-\beta|\left|f\left(x_{0}\right)\right|\left|f\left(x_{1}\right)-f\left(x_{0}\right)\right| \\
& \quad \leq \frac{C_{1}}{h_{n}} \sum_{x \in \mathbb{S}_{n}^{-}}\left(f\left(x^{+}\right)-f(x)\right)^{2}+C_{3}\left|f\left(x_{0}\right)\right|\left|f\left(x_{1}\right)-f\left(x_{0}\right)\right| \\
& \quad \leq \frac{C_{1}}{h_{n}} \sum_{x \in \mathbb{S}_{n}^{-}}\left(f\left(x^{+}\right)-f(x)\right)^{2}+C_{4}\left(f\left(x_{0}\right)^{2} h_{n}+\frac{\left(f\left(x_{0}\right)-f\left(x_{1}\right)\right)^{2}}{h_{n}}\right) \\
& \quad \leq \frac{C_{5}}{h_{n}} \sum_{x \in \mathbb{S}_{n}^{-}}\left(f\left(x^{+}\right)-f(x)\right)^{2}+C_{4} f\left(x_{0}\right)^{2} h_{n}
\end{aligned}
$$

where the constants $C_{1}, \ldots, C_{5}>0$ are independent of $n$ and $f$. Note that with $f\left(x_{n+1}\right)=0$, the following holds:

$$
\begin{aligned}
\sum_{x \in \mathbb{S}_{n}^{-}}\left(f\left(x^{+}\right)-f(x)\right)^{2}= & \sum_{x \in \mathbb{S}_{n}^{-}} f\left(x^{+}\right)^{2}-2 \sum_{x \in \mathbb{S}_{n}^{-}} f\left(x^{+}\right) f(x)+\sum_{x \in \mathbb{S}_{n}^{-}} f(x)^{2} \\
= & 2 \sum_{x \in \mathbb{S}_{n}^{\circ}} f(x)^{2}-\sum_{x \in \mathbb{S}_{n}^{\circ}} f(x) f\left(x^{-}\right)-\sum_{x \in \mathbb{S}_{n}^{\circ}} f(x) f\left(x^{+}\right) \\
& +f\left(x_{0}\right)^{2}-f\left(x_{0}\right) f\left(x_{1}\right) \\
= & \sum_{x \in \mathbb{S}_{n}^{\circ}} f(x)\left(-f\left(x^{-}\right)+2 f(x)-f\left(x^{+}\right)\right) \\
& +f\left(x_{0}\right)^{2}-f\left(x_{0}\right) f\left(x_{1}\right)
\end{aligned}
$$

Then

$$
\begin{aligned}
& \sum_{x \in \mathbb{S}_{n}^{-}} \frac{\delta^{+} x}{s_{n}(x)}\left(\nabla^{+} f(x)\right)^{2}+\frac{1-\beta}{s_{n}\left(x_{0}\right)} f\left(x_{0}\right) \nabla^{+} f\left(x_{0}\right) \\
& \leq \frac{C_{5}}{h_{n}} \sum_{x \in \mathbb{S}_{n}^{o}} f(x)\left(-f\left(x^{-}\right)+2 f(x)-f\left(x^{+}\right)\right)+\frac{C_{5}}{h_{n}} f\left(x_{0}\right)\left(f\left(x_{0}\right)-f\left(x_{1}\right)\right) \\
& \quad \quad+C_{4} f\left(x_{0}\right)^{2} h_{n} \\
& =C_{5} \frac{1}{h_{n}} f_{n}^{T} A f_{n}+C_{4} f\left(x_{0}\right)^{2} h_{n}
\end{aligned}
$$

where $f_{n}=\left(f\left(x_{0}\right), f\left(x_{1}\right), \ldots, f\left(x_{n}\right)\right)^{T} \in \mathbb{R}^{n+1}$, and $A$ is an $(n+1) \times(n+1)$ tridiagonal matrix with diagonal elements 2 (the first diagonal entry is equal to

1) and off-diagonal elements -1. It can easily be seen that

$$
0 \leq \frac{\sum_{x \in \mathbb{S}_{n}^{-}} k(x) f(x)^{2} M_{n}(x)}{(f, f)_{n}} \leq C_{6} \frac{\sum_{x \in \mathbb{S}_{n}^{-}} f(x)^{2} M_{n}(x)}{\sum_{x \in \mathbb{S}_{n}^{-}} f(x)^{2} M_{n}(x)} \leq C_{6}
$$

for some constant $C_{6}>0$ independent of $n$ and $f$, as all terms are positive and $k(x)$ is bounded. Lastly,

$$
\begin{aligned}
(f, f)_{n} & =\sum_{x \in \mathbb{S}_{n}^{-}} f(x)^{2} M_{n}(x) \\
& \geq \sum_{x \in \mathbb{S}_{n}^{\circ}} f(x)^{2} m_{n}(x) \delta x+f\left(x_{0}\right)^{2} M_{n}\left(x_{0}\right) h_{n} \\
& \geq C_{7} h_{n} f_{n}^{T} f_{n}
\end{aligned}
$$

for a constant $C_{7}>0$ independent of $n$ and $f$, as $h_{n} \in(0, \delta)$ with $\delta$ small enough. Hence,

$$
\frac{\left(f,-G_{n} f\right)_{n}}{(f, f)_{n}} \leq \frac{C_{5} \frac{1}{h_{n}} f_{n}^{T} A f_{n}}{C_{7} h_{n} f_{n}^{T} f_{n}}+\frac{C_{4} h_{n} f\left(x_{0}\right)^{2}}{C_{7} h_{n} f_{n}^{T} f_{n}}+C_{6} .
$$

Putting these results into (B.4), one obtains

$$
\lambda_{k}^{n} \leq \frac{C_{5}}{C_{7} h_{n}^{2}} \min _{U_{k}} \max _{f \in U_{k}} \frac{f_{n}^{T} A f_{n}}{f_{n}^{T} f_{n}}+C_{8}
$$

As

$$
\min _{U_{k}} \max _{f \in U_{k}} \frac{f_{n}^{T} A f_{n}}{f_{n}^{T} f_{n}}
$$

is the $k$ th eigenvalue of the matrix $A$, one can use the results of [6, Table 2] or alternatively [7, Theorem 2] to obtain that

$$
\lambda_{k}(A)=4 \sin ^{2} \frac{(2 k-1) \pi}{4 n+6} \leq \frac{4 k^{2} \pi^{2}}{(n+1)^{2}}, \quad k=1,2, \ldots, n+1 .
$$

This now shows that

$$
\lambda_{k}^{n} \leq \frac{C_{5}}{C_{7} h_{n}^{2}} \frac{4 k^{2} \pi^{2}}{(n+1)^{2}}+C_{8} \leq \frac{4 \pi^{2} C_{5}}{C_{7} C_{9}^{2}} k^{2}+C_{8} \leq C_{10} k^{2},
$$

as by Assumption 3, we have $C_{9} \leq h_{n}(n+1)$.

Lemma 3. Consider a grid such that $h_{n} \in(0, \delta)$ with $\delta$ small enough; then there exists a constant $C>0$ such that for any $1 \leq k \leq h_{n}^{-1 / 4}$,

$$
\lambda_{k}^{n}-\lambda_{k} \geq-C h_{n}^{1 / 4}
$$

where $C$ is independent of $k$ and $n$.
Proof. It should first be noted that as $\mathcal{G}$ is a self-adjoint operator, the minmax principle holds (see [3, Theorem 2.1]). In particular,

$$
\lambda_{k}=\min _{L \subset \mathcal{D}, \operatorname{dim} L=k} \max _{\psi \in L, \psi \neq 0} \frac{(\psi,-\mathcal{G} \psi)}{(\psi, \psi)}
$$

where $L$ is a linear subspace of the domain of $\mathcal{G}$.
For $i=1, \ldots, k$ define $\psi_{i}: \mathbb{S} \rightarrow \mathbb{R}$ as a linear interpolation of the approximate eigenfunction $\varphi_{i}^{n}$ over the interval $\mathbb{S}$, which is given by

$$
\psi_{i}(x)=\varphi_{i}^{n}\left(y^{-}\right)+\nabla^{-} \varphi_{i}^{n}(y)\left(x-y^{-}\right)
$$

for $x \in\left[y^{-}, y\right]$ and $y \in \mathbb{S}_{n}^{+}$. Then $\left\{\psi_{1}, \ldots, \psi_{k}\right\}$ form a $k$-dimensional linear space. Furthermore, set $\psi_{a}(x)=\sum_{i=1}^{k} a_{i} \psi_{i}(x)$ where the $a_{i}$ are normalized, i.e. $\sum_{i=1}^{k} a_{i}^{2}=1$. Using the min-max principle and integration by parts, we obtain

$$
\begin{align*}
\lambda_{k} \leq & \max _{a_{1}, \ldots, a_{k}: \sum_{i=1}^{k}} a_{i}^{a_{i}=1} \frac{\left(\psi_{a},-\mathcal{G} \psi_{a}\right)}{\left(\psi_{a}, \psi_{a}\right)} \\
& =\max _{\sum_{i=1}^{k} a_{i}^{2}=1} \frac{\int_{l}^{r} \frac{\psi_{a}^{\prime}(x)^{2}}{s(x)} d x+\int_{l}^{r} k(x) \psi_{a}(x)^{2} M(d x)}{\int_{l}^{r} \psi_{a}(x)^{2} M(d x)} . \tag{B.5}
\end{align*}
$$

We will now estimate the different terms appearing in this equation. First, note that as $1-\beta=0$ for Scheme 1 and $|1-\beta| \leq C \delta^{+} x_{0}$ for Scheme 2, it holds that

$$
\begin{align*}
|1-\beta|\left|\psi_{a}\left(x_{0}\right) \nabla^{+} \psi_{a}\left(x_{0}\right)\right| & \leq C_{1} \psi_{a}\left(x_{0}\right)^{2} \sqrt{\delta^{+} x_{0}}+C_{1}\left(\nabla^{+} \psi_{a}\left(x_{0}\right)\right)^{2}\left(\delta^{+} x_{0}\right)^{3 / 2} \\
& \leq C_{2} \sqrt{h_{n}}+C_{2} \sqrt{\delta} \sum_{x \in \mathbb{S}_{n}^{-}}\left(\nabla^{+} \psi_{a}(x)\right)^{2} \delta^{+} x, \tag{B.6}
\end{align*}
$$

with constants $C_{1}, C_{2}>0$ independent of $a$ and $n$, because
$\psi_{a}\left(x_{0}\right)^{2} \leq \sum_{x \in \mathbb{S}_{n}^{-}} \psi_{a}\left(x_{0}\right) \frac{M_{n}(x)}{M_{n}(x)} \leq C_{3} \sum_{i_{1}=1}^{k} \sum_{i_{2}=1}^{k} a_{i_{1}} a_{i_{2}} \sum_{x \in \mathbb{S}_{n}^{-}} \varphi_{i_{1}}^{n}(x) \varphi_{i_{2}}^{n}(x) M_{n}(x)=C_{3}$,
as

$$
\sum_{x \in \mathbb{S}_{n}^{-}} \varphi_{i_{1}}^{n}(x) \varphi_{i_{2}}^{n}(x) M_{n}(x)=\left(\varphi_{i_{1}}^{n}, \varphi_{i_{2}}^{n}\right)_{n}=\delta_{i_{1}, i_{2}}
$$

and $C_{3}>0$ is independent of $a$ and $n$. Using this result and the fact that $\psi_{a}$ is a piecewise linear function, we obtain

$$
\begin{align*}
\int_{l}^{r} & \frac{1}{s(x)} \psi_{a}^{\prime}(x)^{2} d x-\sum_{x \in \mathbb{S}_{n}^{-}} \frac{1}{s_{n}(x)}\left(\nabla^{+} \psi_{a}(x)\right)^{2} \delta^{+} x-\frac{1-\beta}{s_{n}\left(x_{0}\right)} \psi_{a}\left(x_{0}\right) \nabla^{+} \psi_{a}\left(x_{0}\right) \\
& \leq \sum_{x \in \mathbb{S}_{n}^{-}}\left(\nabla^{+} \psi_{a}(x)\right)^{2} \int_{x}^{x^{+}}\left|\frac{1}{s(y)}-\frac{1}{s_{n}(x)}\right| d y-\frac{1-\beta}{s_{n}\left(x_{0}\right)} \psi_{a}\left(x_{0}\right) \nabla^{+} \psi_{a}\left(x_{0}\right) \\
& \leq C_{4} h_{n} \sum_{x \in \mathbb{S}_{n}^{-}}\left(\nabla^{+} \psi_{a}(x)\right)^{2} \delta^{+} x+C_{5}\left(\sqrt{h_{n}}+\sqrt{\delta} \sum_{x \in \mathbb{S}_{n}^{-}}\left(\nabla^{+} \psi_{a}(x)\right)^{2} \delta^{+} x\right) \\
& \leq C_{6} h_{n} \sum_{x \in \mathbb{S}_{n}^{-}}\left(\nabla^{+} \psi_{a}(x)\right)^{2} \delta^{+} x+C_{5} \sqrt{h_{n}}, \tag{B.7}
\end{align*}
$$

where $C_{4}, C_{5}, C_{6}>0$ are independent of $a$ and $n$. The term appearing in (B.7)
can be handled as follows:

$$
\begin{align*}
\sum_{x \in \mathbb{S}_{n}^{-}} & \left(\nabla^{+} \psi_{a}(x)\right)^{2} \delta^{+} x \\
\leq & -C_{7} \frac{\beta}{s_{n}\left(x_{0}\right)} \psi_{a}\left(x_{0}\right) \nabla^{+} \psi_{a}\left(x_{0}\right)+C_{7} \frac{\beta}{s_{n}\left(x_{0}\right)} \psi_{a}\left(x_{0}\right) \nabla^{+} \psi_{a}\left(x_{0}\right) \\
& +C_{7} \sum_{x \in \mathbb{S}_{n}^{-}} \frac{1}{s_{n}(x)}\left(\nabla^{+} \psi_{a}(x)\right)^{2} \delta^{+} x+C_{7} \sum_{x \in \mathbb{S}_{n}^{-}} k(x) M_{n}(x) \psi_{a}(x)^{2} \\
=- & C_{7} \beta \rho M_{n}\left(x_{0}\right) \psi_{a}\left(x_{0}\right) \nabla^{+} \psi_{a}\left(x_{0}\right)+\frac{C_{7} \beta}{s_{n}\left(x_{0}\right)} \psi_{a}\left(x_{0}\right) \nabla^{+} \psi_{a}\left(x_{0}\right) \\
& \quad-C_{7} \sum_{x \in \mathbb{S}_{n}^{o}} \psi_{a}(x) \delta^{-} x \nabla^{-}\left(\frac{1}{s_{n}(x)} \nabla^{+} \psi_{a}(x)\right)-\frac{C_{7}}{s_{n}\left(x_{0}\right)} \psi_{a}\left(x_{0}\right) \nabla^{+} \psi_{a}\left(x_{0}\right) \\
& +C_{7} \sum_{x \in \mathbb{S}_{n}^{-}} k(x) M_{n}(x) \psi_{a}(x)^{2} \\
=- & C_{7} \sum_{x \in \mathbb{S}_{n}^{-}} \psi_{a}(x) M_{n}(x) G_{n} \psi_{a}(x)+\frac{C_{7}(\beta-1)}{s_{n}\left(x_{0}\right)} \psi_{a}\left(x_{0}\right) \nabla^{+} \psi_{a}\left(x_{0}\right) \\
\leq & C_{7} \sum_{i_{1}=1}^{k} \sum_{i_{2}=1}^{k} a_{i_{1}} a_{i_{2}} \lambda_{i_{2}}^{n} \sum_{x \in \mathbb{S}_{n}^{-}} \varphi_{i_{1}}^{n}(x) \varphi_{i_{2}}^{n}(x) M_{n}(x) \\
& +C_{8}\left(\sqrt{h_{n}}+\sqrt{\delta} \sum_{x \in \mathbb{S}_{n}^{-n}}\left(\nabla^{+} \psi_{a}(x)\right)^{2} \delta^{+} x\right) \\
\leq & C_{9} \lambda_{k}^{n}+C_{7} \sqrt{\delta} \sum_{x \in \mathbb{S}_{n}^{-}}\left(\nabla^{+} \psi_{a}(x)\right)^{2} \delta^{+} x, \tag{B.8}
\end{align*}
$$

where the last inequality follows from the fact that $0 \leq \lambda_{1}^{n}<\lambda_{2}^{n}<\cdots<\lambda_{k}^{n}$. The constants $C_{7}, C_{8}, C_{9}>0$ are independent of $a$ and $n$. We can now choose $\delta$ small enough so that $1-C_{7} \sqrt{\delta}>0$; then

$$
\sum_{x \in \mathbb{S}_{n}^{-}}\left(\nabla^{+} \psi_{a}(x)\right)^{2} \delta^{+} x \leq \frac{C_{9}}{1-C_{7} \sqrt{\delta}} \lambda_{k}^{n} \leq C_{10} \lambda_{k}^{n} .
$$

Combining (B.7) and the previous results yields

$$
\begin{gathered}
\int_{l}^{r} \frac{1}{s(x)} \psi_{a}^{\prime}(x) d x \leq \frac{1-\beta}{s_{n}\left(x_{0}\right)} \psi_{a}\left(x_{0}\right) \nabla^{+} \psi_{a}\left(x_{0}\right)+\sum_{x \in \mathbb{S}_{n}^{-}} \frac{1}{s_{n}(x)}\left(\nabla^{+} \psi_{a}(x)\right)^{2} \delta^{+} x \\
\quad+C_{11} \lambda_{k}^{n} h_{n}+C_{5} h_{n}^{1 / 2} .
\end{gathered}
$$

The second term in (B.5) can be handled as follows:

$$
\begin{aligned}
& \left|\int_{l}^{r} k(x) \psi_{a}(x)^{2} M(d x)-\sum_{x \in \mathbb{S}_{n}^{-}} k(x) \psi_{a}(x)^{2} M_{n}(x)\right| \\
& \quad \leq\left|k\left(x_{0}\right) \psi_{a}\left(x_{0}\right)^{2} M\left(x_{0}\right)-k\left(x_{0}\right) \psi_{a}\left(x_{0}\right)^{2} M_{n}\left(x_{0}\right)\right| \\
& \quad+\left|\int_{l}^{r} k(x) \psi_{a}(x)^{2} m(x) d x-\sum_{x \in \mathbb{S}_{n}^{\circ}} k(x) \psi_{a}(x)^{2} m_{n}(x) \delta x\right| \\
& \quad \leq C_{12} h_{n} \psi_{a}\left(x_{0}\right)^{2}+\frac{1}{2} \sum_{x \in \mathbb{S}_{n}^{-}} \int_{x}^{x^{+}}\left|k(y) \psi_{a}(y)^{2} m(y)-k(x) \psi_{a}(x)^{2} m_{n}(x)\right| d y \\
& \quad+\frac{1}{2} \sum_{x \in \mathbb{S}_{n}^{-}} \int_{x}^{x^{+}}\left|k(y) \psi_{a}(y)^{2} m(y)-k\left(x^{+}\right) \psi_{a}\left(x^{+}\right)^{2} m_{n}\left(x^{+}\right)\right| d y \\
& \leq
\end{aligned}
$$

where the last inequality follows in the same way as in the proof of [9, Lemma 4]. Using this result to bound the numerator, one obtains

$$
\begin{aligned}
& \int_{l}^{r} \frac{1}{s(x)} \psi_{a}^{\prime}(x)^{2} d x+\int_{l}^{r} k(x) \psi_{a}(x)^{2} M(d x) \\
& \quad \leq-\sum_{x \in \mathbb{S}_{n}^{n}} \psi_{a}(x) M_{n}(x) G_{n} \psi_{a}(x)+C_{14}\left(\left(\sqrt{\lambda_{k}^{n}}+\lambda_{k}^{n}\right) h_{n}+\lambda_{k}^{n} h_{n}^{2}+h_{n}^{1 / 2}\right) \\
& \quad \leq \lambda_{k}^{n}+C_{14}\left(\left(\sqrt{\lambda_{k}^{n}}+\lambda_{k}^{n}\right) h_{n}+\lambda_{k}^{n} h_{n}^{2}+h_{n}^{1 / 2}\right)
\end{aligned}
$$

for some constant $C_{14}>0$ independent of $a, k$, and $n$. The denominator can be estimated similarly as before by setting $k(x)=1$ :

$$
\left|\int_{l}^{r} \psi_{a}(x)^{2} M(d x)-1\right| \leq C_{15}\left(\sqrt{\lambda_{k}^{n}} h_{n}+\lambda_{k}^{n} h_{n}^{2}\right)
$$

where $C_{15}>0$ is a constant independent of $a, k$ and $n$. As all of the constants are independent of $a$, it follows that

$$
\lambda_{k} \leq \frac{\lambda_{k}^{n}+C_{14}\left(\left(\sqrt{\lambda_{k}^{n}}+\lambda_{k}^{n}\right) h_{n}+\lambda_{k}^{n} h_{n}^{2}+h_{n}^{1 / 2}\right)}{1-C_{15}\left(\sqrt{\lambda_{k}^{n}} h_{n}+\lambda_{k}^{n} h_{n}^{2}\right)} .
$$

Using Lemma 2, i.e., $\lambda_{k}^{n} \leq C_{16} k^{2} \leq C_{16} h_{n}^{-1 / 2}$ for some $C_{16}>0$ independent of
$k$ and $n$, we have

$$
\begin{aligned}
\lambda_{k}-\lambda_{k}^{n} & \leq \frac{C_{14}\left(\left(\sqrt{\lambda_{k}^{n}}+\lambda_{k}^{n}\right) h_{n}+\lambda_{k}^{n} h_{n}^{2}+h_{n}^{1 / 2}\right)+C_{15} \lambda_{k}^{n}\left(\sqrt{\lambda_{k}^{n}} h_{n}+\lambda_{k}^{n} h_{n}^{2}\right)}{1-C_{15}\left(\sqrt{\lambda_{k}^{n}} h_{n}+\lambda_{k}^{n} h_{n}^{2}\right)} \\
& \leq \frac{C_{17}\left(h_{n}^{3 / 4}+h_{n}^{1 / 2}+h_{n}^{3 / 2}+h_{n}^{1 / 4}+h_{n}^{1 / 2}\right)}{1-C_{15}\left(\delta^{3 / 4}+\delta^{3 / 2}\right)} \leq C_{18} h_{n}^{1 / 4}
\end{aligned}
$$

for constants $C_{17}, C_{18}>0$ independent of $k$ and $n$, as long as $\delta$ is small enough so that $1-C_{15}\left(\delta^{3 / 4}+\delta^{3 / 2}\right)>0$.

Lemma 4. If $h_{n} \in(0, \delta)$ for $\delta$ small enough, there exists a constant $C>0$ such that for any $1 \leq k \leq n$,

$$
\left\|\varphi_{k}^{n}\right\|_{n, \infty} \leq C k
$$

where $C$ is independent of $k$ and $n$.
Proof. Note that for every $y \in \mathbb{S}_{n}^{-}$,

$$
\varphi_{k}^{n}(y)=\sum_{y \leq x<x_{n+1}} \varphi_{k}^{n}(x)-\varphi_{k}^{n}\left(x^{+}\right)=-\sum_{y \leq x<x_{n+1}} \nabla^{+} \varphi_{k}^{n}(x) \delta^{+} x,
$$

as $\varphi_{k}^{n}\left(x_{n+1}\right)=0$. Then

$$
\begin{aligned}
\left|\varphi_{k}^{n}(y)\right| & =\left|-\sum_{y \leq x<x_{n+1}} \nabla^{+} \varphi_{k}^{n}(x) \delta^{+} x\right| \leq \sum_{x \in \mathbb{S}_{n}^{-}}\left|\nabla^{+} \varphi_{k}^{n}(x)\right| \delta^{+} x \\
& \leq \sqrt{\sum_{x \in \mathbb{S}_{n}^{-}}\left(\nabla^{+} \varphi_{k}^{n}(x)\right)^{2} \delta^{+} x \sum_{x \in \mathbb{S}_{n}^{-}} \delta^{+} x} \leq C_{1} \sqrt{\lambda_{k}^{n}}
\end{aligned}
$$

for a constant $C_{1}>0$ independent of $k, n$, and $y$, because of the same steps as shown in (B.8) with $\psi_{a}$ replaced by $\varphi_{k}^{n}$. Furthermore, by Lemma 2, i.e., $\lambda_{k}^{n} \leq C_{2} k^{2}$, it follows that

$$
\left\|\varphi_{k}^{n}\right\|_{n, \infty} \leq C_{1} \sqrt{\lambda_{k}^{n}} \leq C_{3} k
$$

for constants $C_{2}, C_{3}>0$ independent of $k$ and $n$.
Lemma 5. It holds that

$$
\left|\sum_{x \in \mathbb{S}_{n}^{-}} f(x) M_{n}(x)-\int_{x_{0}}^{x_{n+1}} f(x) M(d x)\right| \leq C \max \left\{\|f\|_{\infty},\left\|f^{\prime \prime}\right\|_{\infty}\right\} h_{n}^{\gamma}
$$

for some constant $C>0$ independent of $n$ and $f$.
Proof. We can prove the following by using Proposition 2 and the trapezoidal rule:

$$
\begin{array}{rl}
\sum_{x \in \mathbb{S}_{n}^{-}} f & f(x) M_{n}(x)-\int_{x_{0}}^{x_{n+1}} f(x) M(d x) \\
= & f\left(x_{0}\right) M_{n}\left(x_{0}\right)-f\left(x_{0}\right) M\left(x_{0}\right)+\sum_{x \in \mathbb{S}_{n}^{\circ}} f(x) m_{n}(x) \delta x \\
& \quad-\int_{x_{0}}^{x_{n+1}} f(x) m(x) d x \\
= & f\left(x_{0}\right) \frac{\delta^{+} x_{0}}{\sigma^{2}\left(x_{0}\right)}\left(M\left(x_{0}\right) \alpha-1\right)+O\left(h_{n}^{2}\right)+\sum_{x \in \mathbb{S}_{n}^{\circ}} f(x) m_{n}(x) \delta x \\
& -\int_{x_{0}}^{x_{n+1}} f(x) m(x) d x \\
\leq & C_{1}\left|M\left(x_{0}\right) \alpha-1\right|\|f\|_{\infty} h_{n}+C_{2}\left\|f^{\prime \prime}\right\|_{\infty} h_{n}^{2} \\
\leq & \begin{cases}C_{3}\|f\|_{\infty} h_{n} \quad \text { for } \alpha=\mu\left(x_{0}\right), \\
\left.C_{2}\left\|f^{\prime \prime}\right\|_{\infty} h_{n}^{2} \quad \text { for } \alpha=\rho \quad \text { (as } M\left(x_{0}\right) \alpha=\frac{1}{\rho} \times \rho=1\right) \\
\leq & C_{4} \max \left\{\|f\|_{\infty},\left\|f^{\prime \prime}\right\|_{\infty}\right\} h_{n}^{\gamma},\end{cases}
\end{array}
$$

where $C_{1}, \ldots, C_{4}>0$ are independent of $n$ and $f$.
Corollary 2. For $h_{n} \in(0, \delta)$, the following lower bound holds for every $1 \leq$ $k \leq h_{n}^{-1 / 4}$ :

$$
\lambda_{k}^{n} \geq C k^{2}
$$

if $\delta$ is sufficiently small and $C>0$ is a constant independent of $k$ and $n$.
Proof. The proof is the same as the proof of [5, Corollary 3.7], using Proposition 3 and Lemma 2.

## C. Pseudocode

We provide pseudocode for our CTMC approximation algorithm for pricing of some payoff $f$ under a diffusion model with sticky lower boundary. Symbols
in bold are vectors or matrices. Round brackets show the index or range of indices (indices start at 0 ), and square brackets create a vector.

```
Algorithm 1: Pricing of some payoff \(f\) under a diffusion model with sticky
lower boundary
1 def CTMCPricing ( \(l, r, \mu, \sigma, \rho, k, x, f, T, n\), scheme):
Data: Parameters of underlying diffusion and product to be priced;
``` scheme indicates the scheme to be used.

Result: Price of the product with payoff \(f\).
2 if \(r==\infty\) :
\(3 \quad r \leftarrow r<\infty\)
\(4 \quad \boldsymbol{x}=\left(x_{0}, \ldots, x_{n+1}\right)=\operatorname{Linspace}(l, r, n+2)\)
\(\mathbf{5} \quad \boldsymbol{\delta}^{+} \leftarrow \boldsymbol{x}(2: n)-\boldsymbol{x}(1: n-1), \boldsymbol{\delta}^{-} \leftarrow \boldsymbol{x}(1: n-1)-\boldsymbol{x}(0: n-2), \boldsymbol{\delta} \leftarrow\)
\[
\frac{1}{2}\left(\boldsymbol{\delta}^{+}+\boldsymbol{\delta}^{-}\right)
\]
\(\boldsymbol{v}_{l} \leftarrow-\frac{\mu(\boldsymbol{x}(1: n-1)) * \boldsymbol{\delta}^{+}}{2 \boldsymbol{\delta}^{-} * \boldsymbol{\delta}}+\frac{\sigma^{2}(\boldsymbol{x}(1: n-1))}{2 \boldsymbol{\delta}^{-} * \boldsymbol{\delta}}, \boldsymbol{v}_{u} \leftarrow \frac{\mu(\boldsymbol{x}(1: n-1)) * \boldsymbol{\delta}^{+}}{2 \boldsymbol{\delta}^{-} * \boldsymbol{\delta}}+\frac{\sigma^{2}(\boldsymbol{x}(1: n-1))}{2 \boldsymbol{\delta}^{-} * \boldsymbol{\delta}}\)
\(\boldsymbol{v}_{d} \leftarrow\left[0,-\boldsymbol{v}_{l}-\boldsymbol{v}_{u}-k(\boldsymbol{x}(1: n-1)), 0\right], \boldsymbol{v}_{u} \leftarrow\left[0, \boldsymbol{v}_{u}(0:\right.\)
\(n-2), 0], \boldsymbol{v}_{l} \leftarrow\left[\boldsymbol{v}_{l}, 0\right]\)
\(8 \quad\) if scheme \(==1\) :
\[
\boldsymbol{v}_{u}(0) \leftarrow \rho /(\boldsymbol{x}(1)-\boldsymbol{x}(0)), \boldsymbol{v}_{d}(0) \leftarrow-\boldsymbol{v}_{u}(0)-k(\boldsymbol{x}(0))
\]
else:
\[
\boldsymbol{v}_{u}(0) \leftarrow \frac{\rho}{\boldsymbol{x}(1)-\boldsymbol{x}(0)+\frac{\rho-\mu(\boldsymbol{x}(0))}{\sigma^{2}(\boldsymbol{x}(0))}(\boldsymbol{x}(1)-\boldsymbol{x}(0))^{2}}, \quad \boldsymbol{v}_{d}(0) \leftarrow-\boldsymbol{v}_{u}(0)-k(\boldsymbol{x}(0))
\]
\[
\boldsymbol{G} \leftarrow\left(\begin{array}{ccccc}
\ddots & \ddots & \ddots & & \\
& \boldsymbol{v}_{l} & \boldsymbol{v}_{d} & \boldsymbol{v}_{u} & \\
& & \ddots & \ddots & \ddots
\end{array}\right)
\]
\(\boldsymbol{f} \leftarrow f(\boldsymbol{x})\)
\(\boldsymbol{P}=\) MatrixExponential \((T \boldsymbol{G}) \boldsymbol{f}\)
if \(x \in x\) :
\(i=\operatorname{FindIndex}(x, \boldsymbol{x})\)
return \(P(i)\)
else:
            \(P(\cdot)=\) CubicSplineInterpolation \((\boldsymbol{x}, \boldsymbol{P})\)
            return \(P(x)\)

In line 4, a uniform grid is generated by the function Linspace, but this can be replaced by a non-uniform grid. The operations in line 6 are understood to be componentwise. Line 12 generates the tridiagonal matrix based on the transition rates in \(\boldsymbol{v}_{l}, \boldsymbol{v}_{d}\), and \(\boldsymbol{v}_{u}\), and \(\boldsymbol{f}\) is the vector of payoffs at each grid point. In line 14 the matrix exponential is calculated for \(T \boldsymbol{G}\) and multiplied by the vector \(\boldsymbol{f}\). This is a matrix-vector multiplication, and the result is again an \((n+1)\)-dimensional vector \(\boldsymbol{P}\). The calculation of the matrix exponential can be done using any one of the algorithms described in Section 3.2. The final step is to return the price corresponding to the starting point \(x\). If the starting point is in the grid \(\boldsymbol{x}\), then line 16 returns the index in the vector corresponding to \(x\). In case \(x\) is not on the grid, we apply cubic spline interpolation to the price vector and obtain the price at \(x\) from the cubic spline function.

\section*{References}
[1] Borodin, A. N. and Salminen, P. (2002). Handbook of Brownian Motion-Facts and Formulae, 2nd edn. Birkhäuser, Basel.
[2] Cherny, A. S. (2002). On the uniqueness in law and the pathwise uniqueness for stochastic differential equations. Theory Prob. Appl. 46, 406-419.
[3] Eschwé, D. and Langer, M. (2004). Variational principles for eigenvalues of self-adjoint operator functions. Integral Equat. Operat. Theory 49, 287-321.
[4] Jacod, J. (1979). Calcul Stochastique et Problèmes de Martingales. Springer, Berlin.
[5] Li, L. and Zhang, G. (2018). Error analysis of finite difference and Markov chain approximations for option pricing. Math. Finance 28, 877-919.
[6] Losonczi, L. (1992). Eigenvalues and eigenvectors of some tridiagonal matrices. Acta Math. Hung. 60, 309-322.
[7] Yueh, W.-C. (2005). Eigenvalues of several tridiagonal matrices. Appl. Math. E-Notes 5, 66-74.
[8] Zhang, G. and Li, L. (2019). Analysis of Markov chain approximation for option pricing and hedging: grid design and convergence behavior. Operat. Res. 67, 407-427.
[9] Zhang, G. and Li, L. (2019). Online companion to analysis of Markov chain approximation for option pricing and hedging: grid design and convergence behavior. Available at https://doi.org/10.1287/opre.2018.1791.

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