

**Supplementary material to "On the theory of body motion in  
confined Stokesian fluids"**

Giuseppe Procopio and Massimiliano Giona\*

*Dipartimento di Ingegneria Chimica Materiali Ambiente,  
Sapienza Università di Roma, via Eudossiana 18, Rome 00184, Italy*

## Abstract

This supplementary material reports the Faxén operators of a sphere up to the second-order with Navier-slip boundary conditions and the regular part of the Stokesian multipoles in the semispace domain, useful for constructing the  $[N]$ -matrix for a sphere near a plane wall.

## I. FAXÉN OPERATORS FOR A SPHERE WITH NAVIER SLIP BOUNDARY CONDITIONS

In the following, Faxén operators for a sphere with Navier-slip boundary conditions, evaluated in [2], are reported.

The zeroth-order Faxén operator is

$$\mathcal{F}_{\beta\alpha} = - \left( \frac{1+2\hat{\lambda}}{1+3\hat{\lambda}} \right) \left( \frac{3}{4} R_p + \frac{1}{8} \frac{R_p^3}{(1+2\hat{\lambda})} \Delta_\xi \right) \delta_{\alpha\beta} \quad (1)$$

where  $\Delta_\xi$  is the Laplacian operator acting on the coordinate of the center of the sphere. The first-order Faxén operator reads

$$\begin{aligned} \mathcal{F}_{\alpha\beta\beta_1} = & - \frac{R_p^3}{6(1+5\hat{\lambda})(1+3\hat{\lambda})} \left\{ \left[ (4+20\hat{\lambda}+15\hat{\lambda}^2)\delta_{\alpha\beta}\nabla_{\beta_1} + (1+5\hat{\lambda}+15\hat{\lambda}^2)\delta_{\alpha\beta_1}\nabla_\beta \right] \right. \\ & \left. + \frac{R_p^2}{10} \left[ (4+12\hat{\lambda}-15\hat{\lambda}^2)\Delta_\xi\nabla_{\beta_1}\delta_{\alpha\beta} + (1+3\hat{\lambda}+15\hat{\lambda}^2)\Delta_\xi\nabla_\beta\delta_{\alpha\beta_1} \right] \right\} \end{aligned} \quad (2)$$

by which it is possible to obtain the Faxén operator for the torque on the sphere

$$\mathcal{T}_{\gamma\alpha} = \frac{\varepsilon_{\alpha\gamma_1\gamma} R_p^3 \nabla_{\gamma_1}}{2(1+3\hat{\lambda})} \quad (3)$$

and for the stresses

$$\mathcal{E}_{\alpha\beta\beta_1} = - \frac{\mathcal{F}_{\alpha\beta\beta_1} + \mathcal{F}_{\alpha\beta_1\beta}}{2} = \left( \frac{5+10\hat{\lambda}}{6+30\hat{\lambda}} + \frac{\Delta_\xi}{12(1+5\hat{\lambda})} \right) \left( \frac{\nabla_\beta\delta_{\alpha\beta_1} + \nabla_{\beta_1}\delta_{\alpha\beta}}{2} \right) \quad (4)$$

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\* massimiliano.giona@uniroma1.it

Finally, the second-order Faxén operator is

$$\begin{aligned} \mathcal{F}_{\alpha\beta\beta_1\beta_2} = & -\frac{R_p^3}{4(1+4\hat{\lambda})(1+7\hat{\lambda})} \left\{ \frac{(1+4\hat{\lambda})(1+7\hat{\lambda})}{1+3\hat{\lambda}} \left[ \delta_{\alpha\beta}\delta_{\beta_1\beta_2} + \hat{\lambda}(\delta_{\beta\beta_1}\delta_{\alpha\beta_2} + \delta_{\beta\beta_2}\delta_{\alpha\beta_1}) \right] \right. \\ & + \frac{R_p^2}{6} \left[ -4\hat{\lambda}^2 \left( \frac{4+21\hat{\lambda}}{1+3\hat{\lambda}} \right) \Delta_\xi \delta_{\beta_1\beta_2} \delta_{\alpha\beta} + 5(1+6\hat{\lambda}) \nabla_{\beta_1\beta_2} \delta_{\alpha\beta} + (1+6\hat{\lambda}+28\hat{\lambda}^2) (\nabla_{\beta\beta_1} \delta_{\alpha\beta_2} + \nabla_{\beta\beta_2} \delta_{\alpha\beta_1}) \right. \\ & \left. (1+12\hat{\lambda}+56\hat{\lambda}^2)(\delta_{\alpha\beta_1}\delta_{\beta\beta_2} + \delta_{\alpha\beta_2}\delta_{\beta\beta_1}) \Delta_\xi \right] \\ & \left. + \frac{R_p^4}{84} \left[ (5+20\hat{\lambda}-56\hat{\lambda}^2) \nabla_{\beta_1\beta_2} \delta_{\alpha\beta} + (1+4\hat{\lambda}+28\hat{\lambda}^2) (\nabla_{\beta\beta_1} \delta_{\alpha\beta_2} + \nabla_{\beta\beta_2} \delta_{\alpha\beta_1}) \right] \Delta_\xi \right\} \end{aligned} \quad (5)$$

## II. REGULAR PART OF THE MULTipoles OF THE GREEN FUNCTION FOR A FLUID BOUNDED BY A PLANE WALL

In this Section, the regular part of multipoles of the Stokes flow, defined in the semispace with no-slip boundary conditions and evaluated at the center of a sphere with radius  $R_p$ , at distance  $h$  from the plane boundary, are reported.

The single-pole is obtained by evaluating the regular part of the Green function with both field  $\mathbf{x}$  and pole point  $\boldsymbol{\xi}$  at the position of the center of the sphere  $(0, 0, h)$ . The regular part of the Green function bounded by a plane wall, expressed in the invariant form [1], reads

$$W_{a\alpha}(\mathbf{x}, \boldsymbol{\xi}) = S_{a\alpha}(\mathbf{x} - \boldsymbol{\xi}_r) - (\boldsymbol{\xi}_r - \boldsymbol{\xi}) \cdot \mathbf{n} J_{\alpha\beta'} [\nabla_{\beta'} S_{a3}(\mathbf{x} - \boldsymbol{\xi}_r) - \frac{(\boldsymbol{\xi}_r - \boldsymbol{\xi}) \cdot \mathbf{n}}{2} \Delta_{\boldsymbol{\xi}_r} S_{a\beta'}(\mathbf{x} - \boldsymbol{\xi}_r)] \quad (6)$$

where  $\mathbf{J} = I - 2\mathbf{n} \otimes \mathbf{n}$  is the mirror operator and  $\mathbf{n}$  is the unit vector normal to the plane wall inward into the fluid, and  $\boldsymbol{\xi}_r = \boldsymbol{\xi} - 2h\mathbf{n}$  is the reflection point by the plane of the pole  $\boldsymbol{\xi}$ . Therefore, the non-vanishing entries are

$$\begin{aligned} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} &= \\ W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} &= -\frac{3}{4h} \\ W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} &= -\frac{3}{2h} \end{aligned} \quad (7)$$

Non-vanishing entries of the gradient of eq. (6) are

$$\begin{aligned}
& \partial_{\xi'_1} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_3} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_3} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{3}{8h^2} \\
& \partial_{\xi'_1} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_1} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_2} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_3} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_1} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_3} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_2} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \frac{3}{8h^2} \\
& \partial_{\xi'_3} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_3} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \frac{3}{4h^2}
\end{aligned}$$

Non-vanishing entries of the Laplacian of eq. (6) are

$$\begin{aligned}
& \Delta_{\xi'} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \Delta_{\xi} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \Delta_{\xi'} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \Delta_{\xi} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \frac{1}{2h^3} \\
& \Delta_{\xi'} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \Delta_{\xi} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \frac{2}{h^3} \tag{8}
\end{aligned}$$

Non-vanishing entries of the second derivatives of eq. (6) performed at the field or at the pole point are

$$\begin{aligned} \partial_{\xi'_1} \partial_{\xi'_1} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_1} \partial_{\xi_1} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_2} \partial_{\xi'_2} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_2} \partial_{\xi_2} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_1} \partial_{\xi'_1} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_1} \partial_{\xi_1} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_2} \partial_{\xi'_2} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_2} \partial_{\xi_2} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_1} \partial_{\xi'_1} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_1} \partial_{\xi_1} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_2} \partial_{\xi'_2} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_2} \partial_{\xi_2} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_3} \partial_{\xi'_3} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_3} \partial_{\xi_3} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_3} \partial_{\xi'_3} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_3} \partial_{\xi_3} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_3} \partial_{\xi'_3} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_3} \partial_{\xi_3} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_1} \partial_{\xi'_2} W_{12}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_1} \partial_{\xi_2} W_{12}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_1} \partial_{\xi'_2} W_{21}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_1} \partial_{\xi_2} W_{21}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_2} \partial_{\xi'_1} W_{12}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_2} \partial_{\xi_1} W_{12}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_2} \partial_{\xi'_1} W_{21}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_2} \partial_{\xi_1} W_{21}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_1} \partial_{\xi'_3} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_1} \partial_{\xi_3} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \\ \partial_{\xi'_1} \partial_{\xi'_3} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) &= \partial_{\xi_1} \partial_{\xi_3} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \end{aligned}$$

$$\begin{aligned}
& \partial_{\xi'_2} \partial_{\xi'_3} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_3} \partial_{\xi_2} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} \partial_{\xi'_3} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_3} \partial_{\xi_2} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{1}{2h^3} \\
& \partial_{\xi'_3} \partial_{\xi'_1} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_1} \partial_{\xi_3} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_3} \partial_{\xi'_1} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_1} \partial_{\xi_3} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_3} \partial_{\xi'_2} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_2} \partial_{\xi_3} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_3} \partial_{\xi'_2} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_2} \partial_{\xi_3} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \frac{1}{4h^3} \tag{9}
\end{aligned}$$

Non-vanishing entries of the second derivatives of eq. (6) performed at the field and pole points are

$$\begin{aligned}
& \partial_{\xi'_1} \partial_{\xi_1} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} \partial_{\xi_2} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{7}{16h^3} \\
& \partial_{\xi'_1} \partial_{\xi_1} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} \partial_{\xi_2} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{5}{16h^3} \\
& \partial_{\xi'_1} \partial_{\xi_1} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} \partial_{\xi_2} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{1}{2h^3} \\
& \partial_{\xi'_3} \partial_{\xi_3} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_3} \partial_{\xi_3} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{5}{4h^3} \\
& \partial_{\xi'_3} \partial_{\xi_3} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{1}{h^3} \\
& \partial_{\xi'_1} \partial_{\xi_2} W_{12}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_1} \partial_{\xi_2} W_{21}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} \partial_{\xi_1} W_{12}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} \partial_{\xi_1} W_{21}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{1}{16h^3} \\
& \partial_{\xi'_1} \partial_{\xi_3} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} =
\end{aligned}$$

$$\begin{aligned}
& \partial_{\xi'_3} \partial_{\xi_1} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} \partial_{\xi_3} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_3} \partial_{\xi_2} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{1}{4h^3} \\
& \partial_{\xi'_1} \partial_{\xi_3} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_3} \partial_{\xi_1} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} \partial_{\xi_3} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_3} \partial_{\xi_2} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \frac{1}{2h^3}
\end{aligned}$$

Non-vanishing entries of the gradient at the field point of the Laplacian of eq. (6) are

$$\begin{aligned}
& \partial_{\xi'_1} \Delta_{\xi'} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_1} \Delta_{\xi} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} \Delta_{\xi'} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_2} \Delta_{\xi} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \frac{15}{8h^4} \\
& \partial_{\xi'_1} \Delta_{\xi'} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_1} \Delta_{\xi} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_2} \Delta_{\xi'} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_2} \Delta_{\xi} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_3} \Delta_{\xi'} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_3} \Delta_{\xi} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi'_3} \Delta_{\xi'} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_3} \Delta_{\xi} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{9}{8h^4} \\
& \partial_{\xi'_3} \Delta_{\xi'} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_3} \Delta_{\xi} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{15}{4h^4}
\end{aligned}$$

Non-vanishing entries of the gradient at the pole point of the Laplacian of eq. (6) are

$$\begin{aligned}
& \partial_{\xi_1} \Delta_{\xi'} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_1} \Delta_{\xi} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi_2} \Delta_{\xi'} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_2} \Delta_{\xi} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{15}{8h^4} \\
& \partial_{\xi_1} \Delta_{\xi'} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_1} \Delta_{\xi} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi_2} \Delta_{\xi'} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_2} \Delta_{\xi} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{3}{8h^4} \\
& \partial_{\xi_3} \Delta_{\xi'} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_3} \Delta_{\xi} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi_3} \Delta_{\xi'} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_3} \Delta_{\xi} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \frac{9}{8h^4} \\
& \partial_{\xi_3} \Delta_{\xi'} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_3} \Delta_{\xi} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{9}{4h^4}
\end{aligned}$$

Non-vanishing entries of the second derivatives of the Laplacian of eq. (6) are

$$\begin{aligned}
& \partial_{\xi_3} \partial_{\xi'_2} \Delta_{\xi'} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_3} \partial_{\xi'_2} \Delta_{\xi} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi_3} \partial_{\xi'_3} \Delta_{\xi'} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_3} \partial_{\xi'_3} \Delta_{\xi} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi_3} \partial_{\xi'_3} \Delta_{\xi'} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi_3} \partial_{\xi'_3} \Delta_{\xi} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \frac{3}{2h^5} \\
& \partial_{\xi_3} \partial_{\xi'_3} \Delta_{\xi'} W_{33}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \frac{6}{h^5}
\end{aligned}$$

Non-vanishing entries of the third order derivatives of eq. (6) performed at the field and pole points are



$$\begin{aligned}
& \partial_{\xi_3} \partial_{\xi'_3} \partial_{\xi'_1} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_3} \partial_{\xi_3} \partial_{\xi_1} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi_3} \partial_{\xi'_3} \partial_{\xi'_2} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_3} \partial_{\xi_3} \partial_{\xi_2} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi_3} \partial_{\xi'_3} \partial_{\xi'_3} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_3} \partial_{\xi_3} \partial_{\xi_3} W_{11}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi_3} \partial_{\xi'_3} \partial_{\xi'_3} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_3} \partial_{\xi_3} \partial_{\xi_3} W_{22}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \frac{3}{4h^4} \\
& \partial_{\xi_3} \partial_{\xi'_3} \partial_{\xi'_1} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_3} \partial_{\xi_3} \partial_{\xi_1} W_{31}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi_3} \partial_{\xi'_3} \partial_{\xi'_2} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_3} \partial_{\xi_3} \partial_{\xi_2} W_{32}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi_3} \partial_{\xi'_3} \partial_{\xi'_1} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_3} \partial_{\xi_3} \partial_{\xi_1} W_{13}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \\
& \partial_{\xi_3} \partial_{\xi'_3} \partial_{\xi'_2} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = \partial_{\xi'_3} \partial_{\xi_3} \partial_{\xi_2} W_{23}(\boldsymbol{\xi}', \boldsymbol{\xi}) \Big|_{\boldsymbol{\xi}'=\boldsymbol{\xi}=(0,0,h)} = -\frac{3}{4h^4}
\end{aligned}$$


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- [1] PROCOPIO, G. & GIONA, M. 2023 Bitensorial formulation of the singularity method for stokes flows. *Maths Engng* **5** (2), 1–34.
- [2] PROCOPIO, G. & GIONA, M. 2024 On the hinch–kim dualism between singularity and faxén operators in the hydromechanics of arbitrary bodies in stokes flows. *Phys. Fluids* **36** (3), 032016.