

# Supplementary Material: Guide for Employing Inhomogeneity-Dependent Constitutive Closures for the Fluid-Particle Drag Force

Gregory Rubinstein, Ali Ozel, Xiaolong Yin, Jos Derksen, Sankaran Sundaresan

## 1 Introduction

This Supplementary Material provides a detailed guide for employing the inhomogeneity-dependent drag models that were developed over the course of this study as constitutive relations for larger-scale simulations of fluidized beds. The extent of inhomogeneities is quantified by one of two sub-grid quantities: the scalar variance of the particle volume fraction and a normalized measure for the drift flux. Each of these sub-filter quantities is estimated using a scale-similar approach. The scalar variance of the particle volume fraction,  $\overline{(\phi')^2}$ , is defined as:

$$\overline{(\phi')^2} = \overline{\phi^2} - \overline{\phi}^2. \quad (1)$$

The drift flux,  $\overline{\phi} \tilde{\mathbf{v}}_{drift}$ , is defined as:

$$\overline{\phi} \tilde{\mathbf{v}}_{drift} = \overline{\mathbf{u}_{slip} \phi} - \tilde{\mathbf{u}}_{slip} \overline{\phi}, \quad (2)$$

A normalized measure for the drift flux is given as:

$$v_d = \frac{(\overline{\phi} \tilde{\mathbf{v}}_{drift}) \cdot \tilde{\mathbf{u}}_{slip}}{\tilde{\mathbf{u}}_{slip} \cdot \tilde{\mathbf{u}}_{slip}}. \quad (3)$$

Since the constitutive closures for the fluid-particle drag force derived in the current study are applied to larger-scale simulations of fluidized beds at the scale of the base fluid grid,  $\overline{\phi}$  is defined as the particle volume fraction at the base grid scale, and  $\tilde{\mathbf{u}}_{slip}$  is the slip velocity at the base grid scale. In terms of these quantities, the dimensionless drag force,  $F$ , is defined as:

$$F = \frac{\overline{\mathbf{F}}_{fp} \cdot \tilde{\mathbf{u}}_{slip}}{3\pi\mu_f d_p \tilde{\mathbf{u}}_{slip} \cdot \tilde{\mathbf{u}}_{slip}}, \quad (4)$$

where  $\mu_f$  is the fluid viscosity,  $d_p$  is the particle diameter, and  $\overline{\mathbf{F}}_{fp}$  is the fluid-particle interaction force per particle.

Section 2 in this document focuses on the  $\overline{(\phi')^2}$ -dependent drag model, and section 3 details the  $v_d$ -dependent drag model.

## 2 $\overline{(\phi')^2}$ -dependent drag model

The  $\overline{(\phi')^2}$ -dependent drag model is proposed as an interpolation of the drag at the high and low  $St$  limits:

$$F(\overline{\phi}, \tilde{St}, \overline{(\phi')^2}) = \alpha(\tilde{St}) F_{highSt}(\overline{\phi}, \overline{(\phi')^2}) + (1 - \alpha(\tilde{St})) F_{lowSt}(\overline{\phi}, \overline{(\phi')^2}), \quad (5)$$

where the interpolation function,  $\alpha$ , is found to be:

$$\alpha(\tilde{St}) = \frac{1}{2} \left( 1 + \frac{\tilde{St} - 7}{\tilde{St} + 7} \right). \quad (6)$$

In equations 5 and 6, the modified Stokes number,  $\tilde{St}$ , is defined as:

$$\tilde{St} = \frac{\rho_p |\tilde{\mathbf{u}}_{slip}| d_p}{18\mu_f (1 - \bar{\phi})}. \quad (7)$$

In equation 5, the high and low  $St$  limit drag relations are written as:

$$F_{highSt}(\bar{\phi}, \overline{(\phi')^2}) = F_{0,highSt}(\bar{\phi}) - a_{highSt}(\bar{\phi}) \frac{\overline{(\phi')^2}/b_{highSt}(\bar{\phi})}{1 + \overline{(\phi')^2}/b_{highSt}(\bar{\phi})}, \quad (8)$$

$$F_{lowSt}(\bar{\phi}, \overline{(\phi')^2}) = F_{0,lowSt}(\bar{\phi}) - a_{lowSt}(\bar{\phi}) \frac{\overline{(\phi')^2}/b_{lowSt}(\bar{\phi})}{1 + \overline{(\phi')^2}/b_{lowSt}(\bar{\phi})}, \quad (9)$$

In equation 8,  $F_{0,highSt}$ ,  $a_{highSt}$ , and  $b_{highSt}$  are defined as:

$$F_{0,highSt}(\bar{\phi}) = 8.54 \frac{\bar{\phi}}{1 - \bar{\phi}} + (1 - \bar{\phi})^3 \left(1 + 4.11\sqrt{\bar{\phi}}\right), \quad (10)$$

$$a_{highSt}(\bar{\phi}) = 5.70 - 221 \left(\frac{\bar{\phi}}{\phi_{max}}\right)^{2.5} \left(1 - \frac{\bar{\phi}}{\phi_{max}}\right)^4, \quad (11)$$

$$b_{highSt}(\bar{\phi}) = 0.0220 - 240 \left(\frac{\bar{\phi}}{\phi_{max}}\right)^6 \left(1 - \frac{\bar{\phi}}{\phi_{max}}\right)^{10}, \quad (12)$$

while in equation 9,  $F_{0,lowSt}$ ,  $a_{lowSt}$ , and  $b_{lowSt}$  are defined as:

$$F_{0,lowSt}(\bar{\phi}) = 6.69 \frac{\bar{\phi}}{1 - \bar{\phi}} + (1 - \bar{\phi})^3 \left(1 + 4.09\sqrt{\bar{\phi}}\right), \quad (13)$$

$$a_{lowSt}(\bar{\phi}) = 9.39 - 2100 \left(\frac{\bar{\phi}}{\phi_{max}}\right)^3 \left(1 - \frac{\bar{\phi}}{\phi_{max}}\right)^7, \quad (14)$$

$$b_{lowSt}(\bar{\phi}) = 0.0220 - 283 \left(\frac{\bar{\phi}}{\phi_{max}}\right)^5 \left(1 - \frac{\bar{\phi}}{\phi_{max}}\right)^{12}. \quad (15)$$

In all of these equations,  $\phi_{max} = 0.64$ .

Using equations 5-15, the  $\overline{(\phi')^2}$ -dependent drag model is defined. However, since  $\overline{(\phi')^2}$  is a sub-grid quantity, a scale-similar approach is needed to estimate it. This methodology is detailed in section 2.1.

## 2.1 Scale-similar approach for estimating $\overline{(\phi')^2}$

Since  $\overline{(\phi')^2}$  is a sub-grid scale measure for the extent of inhomogeneities, we employ a scale-similar approach to estimate this quantity. In order to utilize this approach,  $\overline{(\phi')^2}$  is estimated using the following scale-similar form:

$$\overline{(\phi')^2} = \kappa'_1 f_1 \left(\frac{\bar{\phi}}{\phi_{max}}\right) g'_1 \left(\frac{\Delta_f}{d_p}\right). \quad (16)$$

In equation 16,  $f_1(\bar{\phi}/\phi_{max})$  is defined as:

$$f_1 \left(\frac{\bar{\phi}}{\phi_{max}}\right) = \left(\frac{\bar{\phi}}{\phi_{max}}\right)^{n_1} \left(1 - \frac{\bar{\phi}}{\phi_{max}}\right)^{m_1}, \quad (17)$$

where  $n_1 = 1.37$ ,  $m_1 = 3.00$ , and  $\phi_{max} = 0.64$ . For the purposes of applications to larger-scale systems,  $g'_1(\Delta_f/d_p)$  is defined as:

$$g'_1 \left(\frac{\Delta_f}{d_p}\right) = \left(\frac{\Delta_f}{d_p}\right)^2, \quad (18)$$

where  $\Delta_f$  is the base fluid grid size of the larger-scale simulation system.

The quantity,  $\kappa'_1$ , which scales with the extent of inhomogeneities present in a fluidized system, is computed as:

$$\kappa'_1 = \frac{\widehat{\phi}^2 - \widehat{\phi}^2}{f_1(\widehat{\phi}/\phi_{max})g'_1(\Delta_{f,test}/d_p) - f_1(\overline{\phi}/\phi_{max})g'_1(\Delta_f/d_p)}, \quad (19)$$

where  $\Delta_{f,test}$  is the test filter size. For the purposes of applying the constitutive drag relation to larger-scale systems,  $\Delta_{f,test}$  is chosen to be  $3\Delta_f$ , as this is the minimum test filter size for a simulation with a base fluid grid size of  $\Delta_f$ . With  $\Delta_{f,test} = 3\Delta_f$ , the test-filtered quantities are computed as:

$$\hat{q}_{i,j,k} = \frac{1}{27} \sum_{i'=i-1}^{i+1} \left( \sum_{j'=j-1}^{j+1} \left( \sum_{k'=k-1}^{k+1} q_{i',j',k'} \right) \right), \quad (20)$$

where  $q$  is a generic system quantity. The use of a test filter in this approach takes advantage of the assumption that due to scale-similarity, the value of  $\kappa'_1$  does not change over a range of length scales. Using equations 16-19, the sub-grid quantity  $\overline{(\phi')^2}$  can be estimated at the base fluid grid size over the entire fluidized system. Plugging this estimate of  $\overline{(\phi')^2}$  into equations 8 and 9, the  $\overline{(\phi')^2}$ -dependent drag model can be fully implemented as a constitutive closure for larger-scale simulations of fluidized beds.

### 3 $v_d$ -dependent drag model

The drift flux-dependent drag model is defined as follows:

$$F(\overline{\phi}, \tilde{St}, v_d) = F_0(\overline{\phi}, \alpha(\tilde{St})) - a(\overline{\phi}) \frac{v_d/b(\overline{\phi})}{1 + v_d/b(\overline{\phi})}, \quad (21)$$

where  $F_0(\overline{\phi}, \alpha(\tilde{St}))$  is defined as:

$$F_0(\overline{\phi}, \alpha(\tilde{St})) = \frac{F_{0,highSt}(\overline{\phi}) + F_{0,lowSt}(\overline{\phi})}{2} + \left( \frac{F_{0,highSt}(\overline{\phi}) - F_{0,lowSt}(\overline{\phi})}{2} \right) \frac{e^{10(\alpha(\tilde{St})-0.6)} - 1}{e^{10(\alpha(\tilde{St})-0.6)} + 1}. \quad (22)$$

In equation 22,  $\alpha(\tilde{St})$  is defined in equation 6, and  $F_{0,highSt}(\overline{\phi})$  and  $F_{0,lowSt}(\overline{\phi})$  are defined as:

$$F_{0,highSt}(\overline{\phi}) = 8.41 \frac{\overline{\phi}}{1 - \overline{\phi}} + (1 - \overline{\phi})^3 \left( 1 + 3.45\sqrt{\overline{\phi}} \right), \quad (23)$$

$$F_{0,lowSt}(\overline{\phi}) = 6.57 \frac{\overline{\phi}}{1 - \overline{\phi}} + (1 - \overline{\phi})^3 \left( 1 + 3.39\sqrt{\overline{\phi}} \right), \quad (24)$$

respectively. In equation 21,  $a(\overline{\phi})$  and  $b(\overline{\phi})$  are defined as:

$$a(\overline{\phi}) = 10.95 - 24.82 \left( \frac{\overline{\phi}}{\phi_{max}} \right)^{0.5} \left( 1 - \frac{\overline{\phi}}{\phi_{max}} \right)^2, \quad (25)$$

$$b(\overline{\phi}) = -0.093 - 6460 \left( \frac{\overline{\phi}}{\phi_{max}} \right)^8 \left( 1 - \frac{\overline{\phi}}{\phi_{max}} \right)^7, \quad (26)$$

where  $\phi_{max} = 0.64$ .

Using equations 21-26, the  $v_d$ -dependent drag model is defined. However, since  $v_d$  is a sub-grid quantity, a scale-similar approach is needed to estimate it. This methodology is detailed in section 3.1.

### 3.1 Scale-similar approach for estimating $v_d$

Since  $v_d$  is a sub-grid scale measure for the extent of inhomogeneities, we employ a scale-similar approach to estimate this quantity. In order to utilize this approach,  $v_d$  is estimated using the following scale-similar form:

$$v_d = \kappa'_2 f_2 \left( \frac{\bar{\phi}}{\phi_{max}} \right) g'_2 \left( \frac{\Delta_f}{d_p} \right). \quad (27)$$

In equation 27,  $f_2(\bar{\phi}/\phi_{max})$  is defined as:

$$f_2 \left( \frac{\bar{\phi}}{\phi_{max}} \right) = \left( \frac{\bar{\phi}}{\phi_{max}} \right)^{n_2} \left( 1 - \frac{\bar{\phi}}{\phi_{max}} \right)^{m_2}, \quad (28)$$

where  $n_2 = 1.44$ ,  $m_2 = 1.84$ , and  $\phi_{max} = 0.64$ . For the purposes of applications to larger-scale systems,  $g'_2(\Delta_f/d_p)$  is defined as:

$$g'_2 \left( \frac{\Delta_f}{d_p} \right) = \left( \frac{\Delta_f}{d_p} \right)^2, \quad (29)$$

where  $\Delta_f$  is the base fluid grid size of the larger-scale simulation system.

The quantity,  $\kappa'_2$ , which scales with the extent of inhomogeneities present in a fluidized system, is computed as:

$$\kappa'_2 = \frac{\bar{\phi} |\widehat{\tilde{\mathbf{u}}}_{slip}| - \widehat{\bar{\phi}} |\widehat{\tilde{\mathbf{u}}}_{slip}|}{f_2 \left( \widehat{\bar{\phi}}/\phi_{max} \right) g'_2(\Delta_{f,test}/d_p) |\widehat{\tilde{\mathbf{u}}}_{slip}| - f_2(\bar{\phi}/\phi_{max}) g'_2(\Delta_f/d_p) |\tilde{\mathbf{u}}_{slip}|}, \quad (30)$$

where  $\Delta_{f,test}$  is the test filter size. For the purposes of applying the constitutive drag relation to larger-scale systems,  $\Delta_{f,test}$  is chosen to be  $3\Delta_f$ . With  $\Delta_{f,test} = 3\Delta_f$ , the test-filtered quantities are computed using equation 20. Using equations 27-30, the sub-grid quantity  $v_d$  can be estimated at the base fluid grid size over the entire fluidized system. Plugging this estimate of  $v_d$  into equation 21, the  $v_d$ -dependent drag model can be fully implemented as a constitutive closure for larger-scale simulations of fluidized beds.

## 4 Summary

In this Supplementary Material, we have detailed how to utilize the fluid-particle drag models that were derived in this study as constitutive closures in larger-scale simulations of fluidized beds. In section 2.1, we showed how to employ the  $(\phi')^2$ -dependent drag model, and in section 3, we showed how to employ the  $v_d$ -dependent drag model. Each of these two drag models accounts for the effects of inhomogeneities within a system, which is not possible with prior drag relations.